# Wideband Nested Cross Approximation for Helmholtz Problems 

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no. 536

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# Wideband Nested Cross Approximation for Helmholtz problems* 

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September 27, 2012

In this article, the construction of nested bases approximations to discretizations of integral operators with oscillatory kernels is presented. The new method has log-linear complexity and generalizes the adaptive cross approximation (ACA) method to highfrequency problems. It allows for a continuous and numerically stable transition from low to high frequencies.

## 1 Introduction

In this article, the efficient numerical solution of Helmholtz problems

$$
\begin{align*}
&-\Delta u-\kappa^{2} u=0  \tag{1a}\\
& \text { in } \Omega^{c}  \tag{1b}\\
& u+\alpha \partial_{\nu} u=u_{0} \text { on } \Gamma:=\partial \Omega
\end{align*}
$$

used to model acoustics and electromagnetic scattering will be considered. Herein, $\kappa$ denotes the wave number and $\Omega^{c}:=\mathbb{R}^{3} \backslash \bar{\Omega}$ the exterior domain of the obstacle $\Omega \subset \mathbb{R}^{3}$. The paramter $\alpha$ and the right-hand side $u_{0}$ appearing in the impedance condition (1b) are given. A convenient way to solve exterior problems is the reformulation as an integral equation $[10,13,12]$ over the boundary $\Gamma$ of the scatterer $\Omega$. The Galerkin discretization leads to large-scale fully populated matrices $A \in \mathbb{C}^{M \times N}$,

$$
\begin{equation*}
a_{i j}=\int_{\Gamma} \int_{\Gamma} K(x, y) \varphi_{i}(x) \psi_{j}(y) \mathrm{d} s_{y} \mathrm{~d} s_{x}, \quad i \in I:=\{1, \ldots, M\}, \quad j \in J:=\{1, \ldots, N\} \tag{2}
\end{equation*}
$$

with test and ansatz functions $\varphi_{i}, \psi_{j}$, having supports $X_{i}:=\operatorname{supp} \varphi_{i}$ and $Y_{j}:=\operatorname{supp} \psi_{j}$, respectively. We consider kernel functions $K$ of the form

$$
\begin{equation*}
K(x, y):=f(x, y) \exp (2 \pi \mathrm{i} \kappa|x-y|) \tag{3}
\end{equation*}
$$

with an arbitrary asymptotically smooth (with respect to $x$ and $y$ ) function $f$, i.e., there are constants $c_{\mathrm{as}, 1}, c_{\mathrm{as}, 2}>0$ such that for $\boldsymbol{\alpha}, \boldsymbol{\beta} \in \mathbb{N}^{3}$

$$
\begin{equation*}
\left|\partial_{x}^{\boldsymbol{\alpha}} \partial_{y}^{\boldsymbol{\beta}} f(x, y)\right| \leq c_{\mathrm{as}, 1} c_{\mathrm{as}, 2}^{p} \boldsymbol{\alpha}!\boldsymbol{\beta}!\frac{|f(x, y)|}{|x-y|^{p}}, \quad p:=|\boldsymbol{\alpha}+\boldsymbol{\beta}| . \tag{4}
\end{equation*}
$$

[^0]An example is $K(x, y)=S(x-y)$ used in the single layer ansatz, where $S(x)=\exp (2 \pi \mathrm{i} \kappa|x|) /(4 \pi|x|)$ denotes the fundamental solution. Notice that the double layer potential $K(x, y)=\partial_{\nu_{y}} S(x-y)$ is of the form (3) only if $\Gamma$, i.e. the unit outer normal $\nu$, is sufficiently smooth.
Depending on the application, low or high-frequency problems are to be solved. For low-frequency problems, i.e. for $\kappa \operatorname{diam} \Omega \leq 1$, the treecode algorithm [5] and fast multipole methods (FMM) [30, 22, 21, 23] were introduced to treat $A$ with $\log$-linear complexity. The panel clustering method [28] is directed towards more general kernel functions. All previous methods rely on degenerate approximations

$$
\begin{equation*}
K(x, y) \approx \sum_{i=1}^{k} u_{i}(x) v_{i}(y), \quad x \in X, y \in Y \tag{5}
\end{equation*}
$$

using a short sum of products of functions $u_{i}$ and $v_{i}$ depending on only one of the two variables $x$ and $y$ chosen from a pair of domains $X \times Y$ which satisfies the far-field condition

$$
\begin{equation*}
\eta_{\text {low }} \operatorname{dist}(X, Y) \geq \max \{\operatorname{diam} X, \operatorname{diam} Y\} \tag{6}
\end{equation*}
$$

with a given parameter $\eta_{\text {low }}>0$. Since replacing the kernel function $K$ in the integrals (2) with degenerate approximations (5) leads to matrices of low rank, a more direct approach to the efficient treatment of matrices (2) are algebraic methods such as mosaic-skeletons [33] and hierarchical matrices [24, 25]. An efficient and comfortable way to construct low-rank approximations is the adaptive cross approximation (ACA) method [6]. The advantage of this approach compared with explicit kernel approximation is that significantly better approximations can be expected due the quasi-optimal approximation properties; cf. [7]. Furthermore, ACA has the practical advantage that only few of the original entries of $A$ are used for its approximation. A second class are wavelet compression techniques [1], which lead to sparse and asymptotically well-conditioned approximations of the coefficient matrix.
It is known that the fundamental solution $S$ (and its derivatives) of any elliptic operator allows for a degenerate approximation (5) on a pair of domains $(X, Y)$ satisfying (6); see [7]. This applies to the Yukawa operator $-\Delta+\kappa^{2}$ for any $\kappa$, because the decay of $S$ benefits from the positive shift $\kappa^{2}$. However, the negative shift $-\kappa^{2}$ in the Helmholtz operator introduces oscillations in $S$. Hence, for high-frequency Helmholtz problems, i.e. for $\kappa \operatorname{diam} \Omega>1$, the wave number $\kappa$ enters the degree of degeneracy $k$ in (5) in a way that $k$ grows linearly with $\kappa$. In addition to this difficulty, the mesh width $h$ of the discretization has to be chosen such that $\kappa h \sim 1$ for a sufficient accuracy of the solution. We assume that

$$
c_{\kappa}:=\kappa h<\frac{1}{4}, \quad h:=\max \left\{\operatorname{diam} X_{i}, \operatorname{diam} Y_{j}, i \in I, j \in J\right\} \sim 1 / \sqrt{N},
$$

which implies that $\kappa \sim \sqrt{N} \sim \sqrt{M}$. Notice that the recent formulation [14] allows to avoid the previous condition and hence leads to significantly smaller $N$. For high-frequency Helmholtz problems, one- and two-level versions [31,32] were presented with complexity $O\left(N^{3 / 2}\right)$ and $O\left(N^{4 / 3}\right)$, respectively. Multi-level algorithms [16, 2] are able to achieve logarithmic-linear complexity. The previous methods are based on an extensive analytic apparatus that is tailored to the kernel function $K$. To overcome the instability of some of the employed expansions at low frequencies, a wideband version of FMM was presented in [15]. The $\mathcal{H}^{2}$-matrix approach presented in [4] is based on the explicit kernel expansions used in $[2,31]$ for two-dimensional problems.
A well-known idea from physical optics is to approximate $K(\cdot, y)$ in a given direction $e \in \mathbb{S}^{2}$ by a plane wave. The desired boundedness of $k$ with respect to $\kappa$ when approximating

$$
\hat{K}(x, y):=K(x, y) \exp (-2 \pi \mathrm{i} \kappa(x-y, e))
$$

can be achieved if (6) is replaced by a condition which depends on $\kappa$ and which is directionally dependent. This is exploited by the fast multipole methods presented in [11, 17, 18, 29]. The aim of this article is to combine this approach with the ease of use of ACA, i.e., our aim is to construct approximations to $A$ with complexity $k N \log N$ using only few of the original entries of $A$. In this sense, this article generalizes ACA (which achieves log-linear complexity only for low-frequencies) to high-frequency Helmholtz problems. An interesting and important property of the new method is that it will allow a continuous and numerically stable transition from low to high wave numbers $\kappa$ by a generalized far-field condition that fades to the usual condition (6) if the wave number becomes small.

The remaining part of this article is organized as follows. In Sect. 2, we prove estimates as (4) for $\hat{K}$ in a cone around $e$. The desired asymptotic smoothness of $\hat{K}$ leads to a far-field condition on the pair of domains ( $X, Y$ ) on which such estimates are valid. In Sect. 3, these conditions will be accounted for by subdividing the matrix indices hierarchically. It will be seen that the number of blocks resulting from this partitioning is too large to allow for hierarchical matrix approximations with log-linear complexity. Therefore, nested bases approximations are required, and in Sect 4 directional $\mathcal{H}^{2}$-matrices will be introduced as a generalization of usual $\mathcal{H}^{2}$-matrices [27] that incorporate a directional hierarchy. Sect. 5 is devoted to the construction of such directional $\mathcal{H}^{2}$-matrices using only few of the entries of $A$. Error estimates for the constructed nested bases are presented and complexity estimates prove the log-linear overall storage and the log-linear number of operations required by the new technique. Finally, Sect. 6 presents numerical experiments that validate our analysis.

## 2 Directional Asymptotic Smoothness

In [7] it is proved that the singularity function of any elliptic second-order partial differential operator is asymptotically smooth. The latter property can be used to prove convergence of ACA and hence the existence of degenerate kernel approximations (5). The wave number $\kappa$ enters the estimates on $k$ in (5) through the expression $c_{\kappa}:=\max \left\{1, \kappa \max _{x \in X, y \in Y}|x-y|\right\}$, which in general becomes unbounded in the limit $\kappa \rightarrow \infty$. For parts $X$ of the domain of small size, i.e. $\kappa \operatorname{diam} X \leq 1$, satisfying (6), $c_{\kappa}$ and hence $k$ in (5) are bounded independently of $\kappa$. This follows from the fact that the recursive construction of domains satisfying (6) ensures that (6) is sharp in the sense that there is a constant $q>1$ such that

$$
\eta_{\text {low }} \operatorname{dist}(X, Y) \leq q \min \{\operatorname{diam} X, \operatorname{diam} Y\}
$$

Hence, if $x \in X$ and $y \in Y$ with $\operatorname{diam} X \approx \operatorname{diam} Y$, then

$$
\max _{x \in X, y \in Y}|x-y| \leq \operatorname{dist}(X, Y)+\operatorname{diam} X+\operatorname{diam} Y \leq\left(\frac{q}{\eta_{\text {low }}}+2\right) \operatorname{diam} X
$$

and thus

$$
\begin{equation*}
c_{\kappa} \leq \max \left\{1,\left(\frac{q}{\eta_{\text {low }}}+2\right) \kappa \operatorname{diam} X\right\} \leq \frac{q}{\eta_{\text {low }}}+2 \tag{7}
\end{equation*}
$$

is bounded independently of $\kappa$. In the other case $\kappa \operatorname{diam} X>1$, we will not be able to prove asymptotic smoothness with bounded constants. However, a similar property can be proved if the far-field condition (6) is replaced with a frequency dependent condition and if the corresponding far field is subdivided into directions.

For the ease of presentation, $K$ defined in (3) will be investigated as a function of $x$ with fixed $y$. Hence, after shifting $x$ to $x+y$ we consider

$$
\hat{K}(x)=f(x) \exp (2 \pi \mathrm{i} \kappa[|x|-(x, e)])
$$

which can be regarded as $K$ divided by the plane wave $\exp (2 \pi \mathrm{i} \kappa(x, e))$ with some given vector $e \in \mathbb{S}^{2}$; cf. Fig. 2. It will be shown that $\hat{K}$ is asymptotically smooth with respect to $x$ in a cone around $e$.


Figure 1: $\operatorname{Re} K\left(x_{1}, x_{2}, 0\right)$ and $\operatorname{Re} \hat{K}\left(x_{1}, x_{2}, 0\right)$ with $e=(0,1,0)^{T}$.
To this end, $\hat{K}$ will be investigated in the coordinates $r:=|x|$ and

$$
\xi:=|x|-(x, e)
$$

As a first step, we investigate the derivative of $\xi$ in direction $e$.
Lemma 1. Let $x \in \mathbb{R}^{3}$ such that $\varphi:=\measuredangle(x, e)>0$ and $2 \sqrt[4]{8} \sin \varphi<1$. Then there is a constant $c_{\mathrm{as}, 3}>0$ such that

$$
\left|\partial_{e}^{p} \xi(x)\right| \leq c_{\mathrm{as}, 3} 2^{p} p!|x|^{1-p}(\sin \varphi)^{2}, \quad p \in \mathbb{N}
$$

Proof. We may assume that $e=e_{1}:=(1,0,0)^{T}$. Hence, $\varphi=\measuredangle\left(x, e_{1}\right)>0$ and $\xi(x)=|x|-x_{1}$. In order to estimate the $p$-th order derivative of $\xi$ with respect to $x_{1}$, we define $b^{2}=|x|^{2}-x_{1}^{2}=(|x| \sin \varphi)^{2}$ and extend $\xi$ regarded as a function in $x_{1}$ to

$$
\hat{\xi}(z):=\sqrt{z^{2}+b^{2}}-z
$$

which is holomorphic in $B_{\rho}\left(x_{1}\right), \rho:=|x| / 2$. Consider $z=\alpha+\mathrm{i} \beta \in B_{\rho}\left(x_{1}\right)$. Then $\alpha>0$ and $\alpha^{2}>\beta^{2}$ due to $\cos \varphi>1 / 2$, and with $A:=\left|z^{2}+b^{2}\right|=\sqrt{B^{2}+4 \alpha^{2} \beta^{2}}, B:=\alpha^{2}-\beta^{2}+b^{2}$, we have

$$
\sqrt{z^{2}+b^{2}}=\sqrt{\frac{1}{2}(A+B)}+i \operatorname{sgn}(\beta) \sqrt{\frac{1}{2}(A-B)}
$$

Due to $\left|\sqrt{x^{2} \pm y^{2}}-|x|\right| \leq y^{2} / \sqrt{x^{2} \pm y^{2}}$ for all $x, y \in \mathbb{R}$, it follows that

$$
\left|\frac{1}{2}(A+B)-\alpha^{2}\right|=\frac{1}{2}\left|\sqrt{\left(\alpha^{2}+\beta^{2}-b^{2}\right)^{2}+4 \alpha^{2} b^{2}}-\left(\alpha^{2}+\beta^{2}-b^{2}\right)\right| \leq 2 \frac{b^{2} \alpha^{2}}{A}
$$

and

$$
\left|\frac{1}{2}(A-B)-\beta^{2}\right|=\frac{1}{2}\left|\sqrt{\left(\alpha^{2}+\beta^{2}+b^{2}\right)^{2}-4 \beta^{2} b^{2}}-\left(\alpha^{2}+\beta^{2}+b^{2}\right)\right| \leq 2 \frac{b^{2} \beta^{2}}{A}
$$

Hence,

$$
\begin{aligned}
\left|\sqrt{z^{2}+b^{2}}-z\right|^{2} & =\left|\sqrt{\frac{1}{2}(A+B)}-\alpha\right|^{2}+\left|\sqrt{\frac{1}{2}(A-B)}-|\beta|^{2} \leq \frac{8 b^{4}}{A^{2}}\left(\frac{\alpha^{4}}{A+B}+\frac{\beta^{4}}{A-B}\right)\right. \\
& \leq \frac{8 b^{4}}{A^{2}}\left(\frac{\alpha^{4}}{2 \alpha^{2}-4 b^{2} \alpha^{2} / A}+\frac{\beta^{4}}{2 \beta^{2}-4 b^{2} \beta^{2} / A}\right)=\frac{4 b^{4}}{A^{2}} \frac{|z|^{2}}{1-2 b^{2} / A} \leq \frac{4 b^{4}}{A} \frac{1}{1-8 \sqrt{2}(\sin \varphi)^{2}}
\end{aligned}
$$

which follows from $|z|^{2} \leq A$ and

$$
\begin{aligned}
A & =\left|z^{2}+b^{2}\right|=\sqrt{|z|^{4}+2 b^{2}\left(\alpha^{2}-\beta^{2}\right)+b^{4}} \geq \sqrt{|z|^{4}+b^{4}} \geq \frac{1}{\sqrt{2}}\left(|z|^{2}+b^{2}\right) \\
& \geq \frac{1}{\sqrt{2}}\left(\left(x_{1}-\rho\right)^{2}+b^{2}\right) \geq \frac{1}{\sqrt{2}}\left(|x|^{2}+\rho^{2}-2 \rho|x|\right)=\frac{|x|^{2}}{4 \sqrt{2}}=\frac{b^{2}}{4 \sqrt{2}(\sin \varphi)^{2}} .
\end{aligned}
$$

From Cauchy's differentiation formula we obtain

$$
\left|\partial_{x_{1}}^{p} \xi(x)\right|=\left|\hat{\xi}^{(p)}\left(x_{1}\right)\right| \leq \frac{p!}{2 \pi} \int_{B_{\rho}\left(x_{1}\right)} \frac{|\hat{\xi}(z)|}{\left|z-x_{1}\right|^{p+1}} \mathrm{~d} z \leq \frac{2^{p} p!}{|x|^{p-1}} \frac{8 \sqrt{2}(\sin \varphi)^{2}}{\sqrt{1-8 \sqrt{2}(\sin \varphi)^{2}}}
$$

Using the previous estimate on the derivatives of $\xi$, we are now able to estimate the derivatives of $\hat{K}$ in direction $e$ and $e_{\perp} \in \mathbb{S}^{2}$ perpendicular to $e$. To this end, we exploit (4) and make use of

$$
\begin{align*}
\left|\partial_{v}^{p} \hat{K}(x)\right| & \leq \sum_{i+j=p}\binom{p}{i}\left|\partial_{v}^{i} \exp (2 \pi \mathrm{i} \kappa \xi)\right|\left|\partial_{v}^{j} f(x)\right| \\
& \leq c_{\mathrm{as}, 1} p!|\hat{K}(x)| \sum_{i+j=p} \frac{\left|\partial_{v}^{i} \exp (2 \pi \mathrm{i} \kappa \xi)\right|}{i!}\left(\frac{c_{\mathrm{as}, 2}}{r}\right)^{j} \tag{8}
\end{align*}
$$

which holds true for any direction $v$. Thus, we require an estimate for $\left|\partial_{v}^{i} \exp (2 \pi \mathrm{i} \kappa \xi)\right|$. This will be done in the following two lemmas for the cases $v=e$ and $v=e_{\perp}$, respectively.

Lemma 2. Let $x$ and $\varphi$ be as in Lemma 1. Let $d>1 / \kappa$ and define $\eta:=\kappa d^{2} / r, \gamma:=\kappa d \sin \varphi$. Then

$$
\left|\partial_{e}^{p} \hat{K}(x)\right| \leq c p!\left(\frac{\rho}{d}\right)^{p}|\hat{K}(x)|, \quad p \in \mathbb{N}
$$

where $c:=2 c_{\mathrm{as}, 1} \exp \left(8 \pi c_{\mathrm{as}, 3} \gamma\right)$ and $\rho:=\max \left\{c_{\mathrm{as}, 2}, 4\right\} \max \{\gamma, \eta\}$.
Proof. In order to estimate $\left|\partial_{e}^{i} \exp (2 \pi \mathrm{i} \kappa \xi)\right|$, we apply Faà di Bruno's formula expressed in terms of Bell polynomials. Using Lemma 1 and $\kappa d>1$, we obtain

$$
\begin{aligned}
\left|\partial_{e}^{i} \exp (2 \pi \mathrm{i} \kappa \xi)\right| & =\left|\sum_{k=0}^{i} \partial_{\xi}^{k} \exp (2 \pi \mathrm{i} \kappa \xi) \sum_{\substack{\sum_{j} j_{\nu}=k \\
\sum_{\nu} \nu j_{\nu}=i}} \frac{i!}{j_{1}!j_{2}!\ldots} \prod_{\ell=1}^{i-k+1}\left(\frac{\partial_{\ell}^{\ell} \xi(x)}{\ell!}\right)^{j_{\ell}}\right| \\
& \leq i!\sum_{k=0}^{i}(2 \pi \kappa)^{k} \sum_{\substack{\sum_{\nu} j_{\nu}=k \\
\sum_{\nu} \nu j_{\nu}=i}} \frac{1}{j_{1}!j_{2}!\ldots} \prod_{\ell=1}^{i-k+1}\left(\frac{c_{\mathrm{as}, 3} 2^{\ell}(\sin \varphi)^{2}}{r^{\ell-1}}\right)^{j_{\ell}}
\end{aligned}
$$

$$
\begin{aligned}
& \leq 2^{i} i!\sum_{k=0}^{i}\left(2 \pi c_{\mathrm{as}, 3} \kappa(\sin \varphi)^{2}\right)^{k} r^{k-i} \sum_{\substack{\sum_{\nu} j_{\nu}=k \\
\sum_{\nu} \nu j_{\nu}=i}} \frac{1}{j_{1}!j_{2}!\ldots} \\
& =\left(\frac{2}{\kappa d^{2}}\right)^{i} i!\sum_{k=0}^{i}\left(2 \pi c_{\mathrm{as}, 3} \gamma^{2}\right)^{k} \eta^{i-k} \sum_{\sum_{\sum_{\nu} j_{\nu}=k}^{\sum_{\nu} \nu j_{\nu}=i}} \frac{1}{j_{1}!j_{2}!\ldots} \\
& <\left(\frac{2 \tilde{\rho}}{d}\right)^{i} i!\sum_{k=0}^{i}\left(2 \pi c_{\mathrm{as}, 3} \gamma\right)^{k} \sum_{\substack{\sum_{\nu} j_{\nu}=k \\
\sum_{\nu} \nu j_{\nu}=i}} \frac{1}{j_{1}!j_{2}!\ldots}
\end{aligned}
$$

where $\tilde{\rho}:=\max \{\eta, \gamma\}$ and $j_{\nu}=0$ for all $\nu>i-k+1$. From the multinomial theorem for $\boldsymbol{j} \in \mathbb{N}^{d}$ and $L:=i-k+1$

$$
\sum_{|\boldsymbol{j}|=k}\binom{k}{\boldsymbol{j}} x^{\boldsymbol{j}}=\left(\sum_{i=1}^{L} x_{i}\right)^{k}
$$

it follows that

$$
\sum_{\substack{\sum_{\nu} j_{\nu}=k \\ \sum_{\nu} \nu j_{\nu}=i}} \frac{1}{j_{1}!\cdots j_{L}!}=\frac{2^{k}}{k!} \sum_{\substack{\sum_{\nu} j_{\nu}=k \\ \sum_{\nu} \nu j_{\nu}=i}}\binom{k}{j} \prod_{\ell=0}^{L}\left(2^{-\ell}\right)^{j_{\ell}} \leq \frac{2^{k}}{k!}\left(\sum_{\ell=0}^{L} 2^{-\ell}\right)^{k} \leq \frac{4^{k}}{k!}
$$

and hence $\left|\partial_{e}^{i} \exp (2 \pi \mathrm{i} \kappa \xi)\right| \leq \tilde{c} i!(2 \tilde{\rho} / d)^{i}$ with $\tilde{c}:=\exp \left(8 \pi c_{\mathrm{as}, 3} \gamma\right)$. Together with (8) this yields

$$
\begin{aligned}
\left|\partial_{e}^{p} \hat{K}(x)\right| & \leq \tilde{c} c_{\mathrm{as}, 1} p!\frac{|\hat{K}(x)|}{d^{p}} \sum_{i+j=p}(2 \tilde{\rho})^{i}\left(\frac{c_{\mathrm{as}, 2} d}{r}\right)^{j} \leq \tilde{c} c_{\mathrm{as}, 1} p!\left(\frac{2 \tilde{\rho}}{d}\right)^{p}|\hat{K}(x)| \sum_{i+j=p}\left(\frac{c_{\mathrm{as}, 2}}{2}\right)^{j} \\
& \leq 2 \tilde{c} c_{\mathrm{as}, 1} p!\left(\frac{\max \left\{c_{\mathrm{as}, 2}, 4\right\} \tilde{\rho}}{d}\right)^{p}|\hat{K}(x)|
\end{aligned}
$$

due to $d / r=\eta /(\kappa d)<\eta \leq \tilde{\rho}$ and $\sum_{j=0}^{p} t^{j} \leq 2 t^{p}$ for $t \geq 2$.
Lemma 3. Let $d, \eta, \gamma$ as in Lemma 2 such that $\eta, \gamma<1$. Then

$$
\left|\partial_{e_{\perp}}^{p} \hat{K}(x)\right| \leq 2 c_{\mathrm{as}, 1} p!\left(\frac{\rho}{d}\right)^{p}|\hat{K}(x)|,
$$

where $\rho:=\max \left\{\frac{12 \pi}{\sqrt{\tilde{\rho}}}, 2 c_{\text {as }, 2}\right\} \tilde{\rho}$ and $\tilde{\rho}:=\max \{\eta, \gamma\}$.
Proof. First, we claim that $\partial_{e_{\perp}}^{i} \exp (2 \pi \mathrm{i} \kappa \xi)$ consists of at most $3^{i}$ summands of the form

$$
g(x):=c_{g} \frac{(2 \pi \mathrm{i} \kappa)^{n}}{r^{n+2 m}}\left(x, e_{\perp}\right)^{2(n+m)-i} \exp (2 \pi \mathrm{i} \kappa \xi)
$$

where $n, m \in \mathbb{N}$ with $2(m+n) \geq i \geq m+n$ and $\left|c_{g}\right| \leq 2^{i} i$ !. This can be seen by induction using

$$
\begin{aligned}
\partial_{e_{\perp}} g(x) & =c_{g} \frac{(2 \pi \mathrm{i} \kappa)^{n+1}}{r^{n+1+2 m}}\left(x, e_{\perp}\right)^{2(n+1+m)-(i+1)} \exp (2 \pi \mathrm{i} \kappa \xi) \\
& +(2 n+2 m-i) c_{g} \frac{(2 \pi \mathrm{i} \kappa)^{n}}{r^{m}}\left(x, e_{\perp}\right)^{2(n+m)-(i+1)} \exp (2 \pi \mathrm{i} \kappa \xi)
\end{aligned}
$$

$$
-(n+2 m) c_{g} \frac{(2 \pi \mathrm{i} \kappa)^{n}}{r^{n+2(m+1)}}\left(x, e_{\perp}\right)^{2(n+m+1)-(i+1)} \exp (2 \pi \mathrm{i} \kappa \xi)
$$

With $\tilde{c}:=4 \pi$ we estimate

$$
\frac{|g(x)|}{i!} \leq \frac{\left|c_{g}\right|}{i!} \frac{(2 \pi \kappa)^{n}}{r^{n+2 m}}\left|\left(x, e_{\perp}\right)\right|^{2(n+m)-i} \leq \tilde{c}^{i} \frac{\kappa^{n}}{r^{n+2 m}}(r \sin \varphi)^{2(n+m)-i} \leq \tilde{c}^{i} \frac{\kappa^{i-(2 m+n)}}{r^{i-n}}\left(\frac{\gamma}{d}\right)^{2(n+m)-i}
$$

Here, we used that $\left|\left(x, e_{\perp}\right)\right| \leq r \sin \varphi=r \gamma /(\kappa d)$. As in the proof of Lemma 2 we have that $d / r<\eta$ and hence

$$
\frac{d^{i}}{i!}|g(x)|=\tilde{c}^{i} \frac{\kappa^{i-(2 m+n)}}{r^{i-n}} \gamma^{2(n+m)-i} d^{2(i-n-m)}=\tilde{c}^{i} \eta^{i-(2 m+n)} \gamma^{2(n+m)-i}\left(\frac{d}{r}\right)^{2 m} \leq \tilde{c}^{i} \tilde{\rho}^{2 m+n} \leq \tilde{c}^{i} \tilde{\rho}^{i / 2}
$$

The last estimate follows from $2 m+n \geq i / 2$. This implies that $\left|\partial_{e_{\perp}}^{i} \exp (2 \pi \mathrm{i} \kappa \xi)\right| \leq i!(3 \tilde{c} / d)^{i} \tilde{\rho}^{i / 2}$ and together with (8) we get

$$
\begin{aligned}
\frac{d^{p}}{p!}\left|\partial_{e_{\perp}}^{p} \hat{K}(x)\right| & \leq c_{\mathrm{as}, 1}|\hat{K}(x)| \sum_{i+j=p}(3 \tilde{c})^{i} \tilde{\rho}^{i / 2}\left(\frac{c_{\mathrm{as}, 2} d}{r}\right)^{j} \leq c_{\mathrm{as}, 1}\left(c_{\mathrm{as}, 2} \tilde{\rho}\right)^{p}|\hat{K}(x)| \sum_{i+j=p}\left(\frac{3 \tilde{c}}{c_{\mathrm{as}, 2} \sqrt{\tilde{\rho}}}\right)^{i} \\
& \leq c_{\mathrm{as}, 1}\left(c_{\mathrm{as}, 2} \tilde{\rho}\right)^{p}|\hat{K}(x)| \sum_{i=0}^{p}\left(\frac{\hat{c}}{c_{\mathrm{as}, 2}}\right)^{i} \leq 2 c_{\mathrm{as}, 1}(\hat{c} \tilde{\rho})^{p}|\hat{K}(x)|,
\end{aligned}
$$

where $\hat{c}:=\max \left\{\frac{3 \tilde{c}}{\sqrt{\tilde{\rho}}}, 2 c_{\mathrm{as}, 2}\right\}$.
We return to the general case of estimating the derivatives of $\hat{K}(x, y)$ for $x \in X$ and $y \in Y$. The last two lemmata show that the derivatives of $\hat{K}$ can be controlled by the parameters $\eta$, $\gamma$, and $d$. Let $\chi(X)$ denote the Chebyshev center of $X$. Using the angle condition

$$
\begin{equation*}
\sin \measuredangle(\chi(X)-y, e) \leq \frac{\gamma_{\text {high }}}{\kappa \operatorname{diam} X}, \quad y \in Y \tag{9}
\end{equation*}
$$

and the high-frequency far-field condition

$$
\begin{equation*}
\eta_{\text {high }} \operatorname{dist}(X, Y) \geq \kappa(\operatorname{diam} X)^{2} \tag{10}
\end{equation*}
$$

with $0<\gamma_{\text {high }}, \eta_{\text {high }}<1$, we obtain for the choice $d=\operatorname{diam} X$ and $x \mapsto x-y$ that $d>1 / \kappa$ and

$$
\eta=\frac{\kappa d^{2}}{r}=\frac{\kappa(\operatorname{diam} X)^{2}}{|x-y|} \leq \frac{\kappa(\operatorname{diam} X)^{2}}{\operatorname{dist}(X, Y)} \leq \eta_{\text {high }}
$$

The following lemma shows that $\gamma=\kappa \operatorname{diam} X \sin \measuredangle(x-y, e)$ is bounded by $\frac{\gamma_{\text {high }}+\eta_{\text {high }}}{1-\eta_{\text {high }}}$.
Lemma 4. Let $X$ and $Y$ satisfy (9) and (10). Then for $x \in X$ and $y \in Y$

$$
\sin \measuredangle(x-y, e) \leq \frac{\gamma_{\text {high }}+\eta_{\text {high }}}{1-\eta_{\text {high }}} \frac{1}{\kappa \operatorname{diam} X}
$$

Proof. Let $u \times v$ denote the cross product of $u, v \in \mathbb{R}^{3}$. Then

$$
|(x-y) \times e| \leq|x-\chi(X)|+|(\chi(X)-y) \times e| \leq \operatorname{diam} X+|(\chi(X)-y) \times e| .
$$

It follows that

$$
\sin \measuredangle(x-y, e)=\frac{|(x-y) \times e|}{|x-y|} \leq \frac{|(\chi(X)-y) \times e|+\operatorname{diam} X}{|\chi(X)-y|-\operatorname{diam} X} .
$$

Due to $\eta_{\text {high }}|\chi(X)-y| \geq \kappa(\operatorname{diam} X)^{2} \geq \operatorname{diam} X$, we obtain that the denominator of the last expression is bounded from below by $\left(1-\eta_{\text {high }}\right)|\chi(X)-y|$. Hence,

$$
\begin{aligned}
\sin \measuredangle(x-y, e) & \leq \frac{1}{1-\eta_{\text {high }}} \frac{|(\chi(X)-y) \times e|+\operatorname{diam} X}{|\chi(X)-y|} \\
& \leq \frac{1}{1-\eta_{\text {high }}}\left(\frac{\gamma_{\text {high }}}{\kappa \operatorname{diam} X}+\frac{\operatorname{diam} X}{|\chi(X)-y|}\right) \leq \frac{\gamma_{\text {high }}+\eta_{\text {high }}}{1-\eta_{\text {high }}} \frac{1}{\kappa \operatorname{diam} X} .
\end{aligned}
$$

As a consequence of the angle condition (9) and the far-field condition (10) we obtain from Lemma 2 and Lemma 3 for $x \in X$ and $y \in Y$

$$
\begin{equation*}
\max \left\{\left|\partial_{e, x}^{p} \hat{K}(x, y)\right|,\left|\partial_{e_{\perp}, x}^{p} \hat{K}(x, y)\right|\right\} \leq c p!\left(\frac{\rho}{\operatorname{diam} X}\right)^{p}|\hat{K}(x, y)| \tag{11}
\end{equation*}
$$

where $c$ is independent of $\kappa$, the directions $e, e_{\perp} \in \mathbb{S}^{2}$ satisfy $\left(e, e_{\perp}\right)=0$ and $0<\rho<1$ for small enough $\eta_{\text {high }}$ and $\gamma_{\text {high }}$.

## 3 Matrix partitioning

The aim of this section is to partition the set of indices $I \times J, I=\{1, \ldots, N\}$ and $J=\{1, \ldots, M\}$, of the matrix defined in (2) into sub-blocks $t \times s, t \subset I$ and $s \subset J$, such that the associated supports

$$
X_{t}:=\bigcup_{i \in t} X_{i} \quad \text { and } \quad Y_{s}:=\bigcup_{j \in s} Y_{j}
$$

satisfy (10) in the high-frequency case $\kappa \operatorname{diam} X_{t}>1$ and (6) if $\kappa \operatorname{diam} X_{t} \leq 1$.
Before we discuss the matrix partition, let us make some assumptions that are in line with the usual finite element discretization. The first assumption is that the overlap of the sets $X_{i}, i \in I$, is bounded in the sense that there is a constant $\nu>0$ such that

$$
\begin{equation*}
\sum_{i \in t} \mu\left(X_{i}\right) \leq \nu \mu\left(X_{t}\right), \quad t \subset I \tag{12}
\end{equation*}
$$

Furthermore, the surface measure $\mu$ has the property that there is $c_{\Gamma}>0$ such that

$$
\mu(X) \leq c_{\Gamma}(\operatorname{diam} X)^{2}
$$

for all $X \subset \Gamma$. The usual way of constructing hierarchical matrix partitions is based on cluster trees; see $[26,7]$. We assume that a binary cluster tree $T_{I}$ for $I$ is constructed such that there are constants $c_{g}, c_{G}>0$ with

$$
\begin{equation*}
2^{-\ell} / c_{G} \leq \mu\left(X_{t}\right) \quad \text { and } \quad\left(\operatorname{diam} X_{t}\right)^{2} \leq c_{g} 2^{-\ell} \tag{13}
\end{equation*}
$$

for all $t$ from the $\ell$-th level $T_{I}^{(\ell)}$ of $T_{I}$. We will make use of the notation $S_{I}(t)$ for the set of sons of a cluster $t \in T_{I}$. The same properties are also assumed for the sets $Y_{j}, j \in J$. Under these assumptions it follows that the depth $L$ of the cluster trees is of the order $L \sim \log N \sim \log M$.

Using the trees $T_{I}$ and $T_{J}$, a partition $P$ can be constructed as the leaves of a block cluster tree $T_{I \times J}$, where we define $S_{I \times J}(t, s)=\emptyset$ for the set of sons of a block $t \times s$ if $t \times s$ satisfies the low-frequency far-field condition

$$
\begin{equation*}
\kappa \min \left\{\operatorname{diam} X_{t}, \operatorname{diam} Y_{s}\right\} \leq 1 \tag{14a}
\end{equation*}
$$

$$
\begin{equation*}
\eta_{\text {low }} \operatorname{dist}\left(X_{t}, Y_{s}\right) \geq \max \left\{\operatorname{diam} X_{t}, \operatorname{diam} Y_{s}\right\} \tag{14b}
\end{equation*}
$$

or the high-frequency far-field condition

$$
\begin{align*}
\kappa \min \left\{\operatorname{diam} X_{t}, \operatorname{diam} Y_{s}\right\} & >1  \tag{15a}\\
\eta_{\text {high }} \operatorname{dist}\left(X_{t}, Y_{s}\right) & \geq \kappa \max \left\{\left(\operatorname{diam} X_{t}\right)^{2},\left(\operatorname{diam} Y_{s}\right)^{2}\right\} \tag{15b}
\end{align*}
$$

or $\min \{|t|,|s|\} \leq n_{\text {min }}$ with some given constants $\eta_{\text {low }}, \eta_{\text {high }}, n_{\text {min }}>0$. In all other cases, we set $S_{I \times J}(t, s)=S_{I}(t) \times S_{J}(s)$. Notice that for $\kappa=0$ we obtain the usual far-field condition (6). Furthermore, the transition from the low to the high-frequency regime is continuous in the sense that for $\kappa \min \left\{\operatorname{diam} X_{t}, \operatorname{diam} Y_{s}\right\}=1$ the conditions (14b) and (15b) are equivalent with $\eta_{\text {high }}=\eta_{\text {low }}$.

As usual, we partition the set $P$ into admissible and non-admissible blocks

$$
P=P_{\mathrm{adm}} \cup P_{\text {nonadm }}
$$

where each $t \times s \in P_{\text {adm }}$ satisfies (14) or (15) and each $t \times s \in P_{\text {nonadm }}$ is small, i.e. satisfies $\min \{|t|,|s|\} \leq n_{\text {min }}$. In order to distinguish low and high-frequency blocks, we further subdivide

$$
P_{\mathrm{adm}}=P_{\text {low }} \cup P_{\text {high }}
$$

where $P_{\text {low }}:=\{t \times s \in P: t \times s$ satisfies (14) $\}$ and $P_{\text {high }}:=P_{\text {adm }} \backslash P_{\text {low }}$.
The following lemma will be the basis for the complexity analysis of the algorithms presented in this article. Notice that this lemma analyzes the so-called sparsity constant of hierarchical matrix partitions introduced in [20] for the far-field condition (6). Since the lemma states that this constant is unbounded with respect to $\kappa$, usual $\mathcal{H}$-matrices are not able to guarantee logarithmic-linear complexity for the high-frequency far-field condition (15b). Therefore, in the next section a variant of $\mathcal{H}^{2}$-matrices will be introduced.
Lemma 5. Let $t \in T_{I}^{(\ell)}$. The set $\left\{s \in T_{J}: t \times s \in P_{\text {high }}\right\}$ has cardinality $\mathcal{O}\left(2^{-\ell} \kappa^{2}\right)$.
Proof. The assumptions (12) and (13) guarantee that each set

$$
N_{\rho}:=\left\{s \in T_{J}^{(\ell)}: \max _{y \in Y_{s}}\left|\chi\left(X_{t}\right)-y\right| \leq \rho\right\}, \quad \rho>0
$$

contains at most $\nu c_{G} c_{\Gamma} 2^{\ell}(2 \rho)^{2}$ clusters $s$ from the same level $\ell$ in $T_{J}$. This follows from

$$
\begin{equation*}
\left|N_{\rho}\right| 2^{-\ell} / c_{G} \leq \sum_{s \in N_{\rho}} \mu\left(Y_{s}\right) \leq \nu \mu\left(Y_{N_{\rho}}\right) \leq \nu c_{\Gamma}(2 \rho)^{2} \tag{16}
\end{equation*}
$$

Let $s \in T_{J}$ such that $t \times s \in P_{\text {high }}$. Furthermore, let $t^{*}$ and $s^{*}$ be the father clusters of $t$ and $s$, respectively. Suppose that $\max _{y \in Y_{s}}\left|\chi\left(X_{t}\right)-y\right| \geq \rho_{0}$, where

$$
\rho_{0}:=\kappa / \eta_{\text {high }} \max \left\{\left(\operatorname{diam} X_{t^{*}}\right)^{2},\left(\operatorname{diam} Y_{s^{*}}\right)^{2}\right\}+\operatorname{diam} X_{t^{*}}+\operatorname{diam} Y_{s^{*}}
$$

Then

$$
\begin{aligned}
\operatorname{dist}\left(X_{t^{*}}, Y_{s^{*}}\right) & \geq \max _{y \in Y_{s}}\left|\chi\left(X_{t}\right)-y\right|-\operatorname{diam} X_{t^{*}}-\operatorname{diam} Y_{s^{*}} \\
& \geq \kappa / \eta_{\text {high }} \max \left\{\left(\operatorname{diam} X_{t^{*}}\right)^{2},\left(\operatorname{diam} Y_{s^{*}}\right)^{2}\right\}
\end{aligned}
$$

implies that $t^{*} \times s^{*} \in P_{\text {high }}$. Hence, $P_{\text {high }}$ cannot contain $t \times s$, which is a contradiction. It follows that

$$
\max _{y \in Y_{s}}\left|\chi\left(X_{t}\right)-y\right|<\rho_{0} \leq\left(c_{g} 2^{-(\ell-1)}\right)^{1 / 2}\left(2+\left(c_{g} 2^{-(\ell-1)}\right)^{1 / 2} \kappa / \eta_{\text {high }}\right)
$$

From (16) we obtain that

$$
\left|\left\{s \in T_{J}: t \times s \in P_{\mathrm{high}}\right\}\right| \leq 8 \nu c_{\Gamma} c_{g} c_{G}\left(2+\left(c_{g} 2^{-(\ell-1)}\right)^{1 / 2} \kappa / \eta_{\mathrm{high}}\right)^{2} \sim 16 \nu c_{\Gamma} c_{g}^{2} c_{G} 2^{-\ell}\left(\kappa / \eta_{\mathrm{high}}\right)^{2}
$$

## Directional subdivision of high-frequency blocks

In the high-frequency regime, i.e. $\kappa \min \left\{\operatorname{diam} X_{t}, \operatorname{diam} Y_{s}\right\}>1$, the matrix block corresponding to $t \times s \in P_{\text {high }}$ cannot be approximated independently of $\kappa$ unless it is directionally subdivided; see the discussion in Sect. 2. In view of the angle condition (9), we partition the space $\mathbb{R}^{3}$ recursively into a hierarchy of unbounded pyramids. The first subdivision partitions $\mathbb{R}^{3}$ into the 6 pyramids defined by the origin and the faces of the unit cube as the pyramids' bases. In each of the next steps, a pyramid is subdivided by dividing its base perpendicular to a largest side of the base into two equally sized halves. A pyramid $Z$ resulting from $\nu$ subdivisions satisfies

$$
\begin{equation*}
\sin \measuredangle(x, e) \leq 2^{(1-\nu) / 2} \quad \text { for all } x \in Z, \tag{17}
\end{equation*}
$$

where $e \in \mathbb{S}^{2}$ denotes the vector pointing from the origin to the center of the base of $Z$. For a given cluster $t \in T_{I}^{(\ell)}$ from the $\ell$-th level of $T_{I}$ let $\nu_{t}$ be the smallest non-negative integer such that

$$
\begin{equation*}
\nu_{t} \geq 2\left(\nu_{0}+\log _{2} \kappa\right)-\ell+1 \tag{18}
\end{equation*}
$$

where $\nu_{0} \in \mathbb{N}$ is a fixed value which will be specified later on. Denote by $\mathcal{E}(t)$ the set of directions $e \in \mathbb{S}^{2}$ associated with all pyramids $Z_{e}$ after $\nu_{t}$ subdivisions. Then

$$
\begin{equation*}
|\mathcal{E}(t)|=6 \cdot 2^{\nu_{t}} \sim \kappa^{2} 2^{-\ell} \tag{19}
\end{equation*}
$$

Given $e \in \mathcal{E}(t)$, we define $e^{\prime} \in \mathcal{E}\left(t^{\prime}\right)$ as the axis of the pyramid $Z_{e^{\prime}}$ from which $Z_{e}$ results after subdivision. Notice that despite $t^{\prime} \subset t$ we have $Z_{e} \subset Z_{e^{\prime}}$. Furthermore, let

$$
\mathcal{F}_{e}\left(X_{t}\right):=\mathcal{D}_{e}\left(X_{t}\right) \cap \mathcal{F}\left(X_{t}\right), \quad \mathcal{D}_{e}(X):=\left\{y \in \mathbb{R}^{3}: \sin \measuredangle(\chi(X)-y, e) \leq \frac{\gamma_{\mathrm{high}}}{\kappa \operatorname{diam} X}\right\}
$$

be the directional far field of $X_{t}$. Here, $\mathcal{F}(X):=\left\{y \in \mathbb{R}^{3}: \eta_{\text {high }} \operatorname{dist}(X, y) \geq \kappa(\operatorname{diam} X)^{2}\right\}$ denotes the far field of $X \subset \mathbb{R}^{3}$.

The following lemma will be important for the construction of nested bases used for the approximation. Notice that the directional far fields are nested up to constants.

Lemma 6. Let $t^{\prime} \in S_{I}(t)$ and $e^{\prime}$ be defined as above. Then the directional far field satisfies

$$
\mathcal{F}_{e}\left(X_{t}\right) \subset \tilde{\mathcal{F}}_{e^{\prime}}\left(X_{t^{\prime}}\right), \quad \tilde{\mathcal{F}}_{e^{\prime}}\left(X_{t^{\prime}}\right):=\left\{y \in \Gamma: \sin \measuredangle\left(\chi\left(X_{t^{\prime}}\right)-y, e^{\prime}\right) \leq \frac{\tilde{\gamma}}{\kappa \operatorname{diam} X_{t^{\prime}}}\right\} \cap \mathcal{F}\left(X_{t^{\prime}}\right)
$$

with the constant $\tilde{\gamma}:=2^{1-\nu_{0}} \sqrt{c_{g}}+\left(\gamma_{\text {high }}+\eta_{\text {high }}\right) /\left(1-\eta_{\text {high }}\right)$.
Proof. For $y \in \mathcal{F}_{e}\left(X_{t}\right)$ let $\zeta=\chi\left(X_{t^{\prime}}\right)-y-\left(\chi\left(X_{t^{\prime}}\right)-y, e\right) e, \zeta^{\prime}:=\chi\left(X_{t^{\prime}}\right)-y-\left(\chi\left(X_{t^{\prime}}\right)-y, e^{\prime}\right) e^{\prime}$, and $\delta:=\zeta^{\prime}-\zeta$. Then

$$
\sin \measuredangle\left(\chi\left(X_{t^{\prime}}\right)-y, e^{\prime}\right)=\frac{\left|\zeta^{\prime}\right|}{\left|\chi\left(X_{t^{\prime}}\right)-y\right|} \leq \frac{|\zeta|+|\delta|}{\left|\chi\left(X_{t^{\prime}}\right)-y\right|}=\sin \measuredangle\left(\chi\left(X_{t^{\prime}}\right)-y, e\right)+\frac{|\delta|}{\left|\chi\left(X_{t^{\prime}}\right)-y\right|}
$$

and

$$
\begin{aligned}
|\delta| & =\left|\left(\chi\left(X_{t^{\prime}}\right)-y, e\right)\left[e-\left(e, e^{\prime}\right) e^{\prime}\right]-\left(\chi\left(X_{t^{\prime}}\right)-y, e^{\prime}-\left(e, e^{\prime}\right) e\right) e^{\prime}\right| \\
& \leq\left|\chi\left(X_{t^{\prime}}\right)-y\right|\left[\left|e-\left(e, e^{\prime}\right) e^{\prime}\right|+\left|e^{\prime}-\left(e, e^{\prime}\right) e\right|\right] \leq 2\left|\chi\left(X_{t^{\prime}}\right)-y\right| \sin \measuredangle\left(e^{\prime}, e\right) .
\end{aligned}
$$

Using Lemma 4, we obtain from $\chi\left(X_{t^{\prime}}\right) \in X_{t^{\prime}} \subset X_{t}$

$$
\sin \measuredangle\left(\chi\left(X_{t^{\prime}}\right)-y, e\right) \leq \frac{\gamma_{\text {high }}+\eta_{\text {high }}}{1-\eta_{\text {high }}} \frac{1}{\kappa \operatorname{diam} X_{t}}
$$

From $e^{\prime} \in \partial Z_{e}$ it follows

$$
\sin \measuredangle\left(e^{\prime}, e\right) \leq 2^{\left(1-\nu_{t}\right) / 2} \leq \frac{2^{-\nu_{0}}}{\kappa} 2^{\ell / 2} \leq \frac{2^{-\nu_{0}} \sqrt{c_{g}}}{\kappa \operatorname{diam} X_{t}}
$$

due to (17) and (13). We obtain

$$
\sin \measuredangle\left(\chi\left(X_{t^{\prime}}\right)-y, e^{\prime}\right) \leq\left(\frac{\gamma_{\text {high }}+\eta_{\text {high }}}{1-\eta_{\text {high }}}+2^{1-\nu_{0}} \sqrt{c_{g}}\right) \frac{1}{\kappa \operatorname{diam} X_{t}} .
$$

The inclusion $\mathcal{F}\left(X_{t}\right) \subset \mathcal{F}\left(X_{t^{\prime}}\right)$ is obvious due to $X_{t^{\prime}} \subset X_{t}$.
From the previous proof it can be seen that $\sin \measuredangle(x, e) \leq \frac{2^{-\nu_{0}} \sqrt{c_{g}}}{\kappa \text { diam } X_{t}}$ for all $x \in Z_{e}$. With Lemma 6 it follows that

$$
\begin{equation*}
\left(\chi\left(X_{t}\right)+Z_{e}\right) \cap \mathcal{F}\left(X_{t}\right) \subset \mathcal{F}_{e^{\prime}}\left(X_{t^{\prime}}\right) \quad \text { for all } e \in \mathcal{E}(t) \tag{20}
\end{equation*}
$$

provided that $\nu_{0}$ from (18) is chosen sufficiently large and $\eta_{\text {high }}$ is chosen sufficiently small.
Since high-frequency blocks require special attention, we gather column clusters $s$ which are admissible with $t$ in direction $e \in \mathcal{E}(t)$ in the (cluster) far field with direction $e$

$$
F_{e}(t):=\bigcup\left\{s \in T_{J}: \exists \hat{t} \supset t \text { such that } \hat{t} \times s \in P_{\text {high }}\right\} \cap\left\{j \in J: Y_{j} \subset \mathcal{F}_{e}\left(X_{t}\right)\right\} .
$$

## 4 Directional $\mathcal{H}^{\mathbf{2}}$-matrices

From Lemma 5 we see that $\mathcal{H}$-matrices are not able to achieve logarithmic linear complexity. Therefore we employ a nested representation similar to $\mathcal{H}^{2}$-matrices introduced in [27]. To account for the required directional subdivision, we generalize the concept of nested cluster bases.

Definition 1. $A$ directional cluster basis $\mathcal{U}$ for the rank distribution $\left(k_{t}^{e}\right)_{t \in T_{I}, e \in \mathcal{E}(t)}$ is a family $\mathcal{U}=$ $\left(U_{e}(t)\right)_{t \in T_{I}, e \in \mathcal{E}(t)}$ of matrices $U_{e}(t) \in \mathbb{C}^{t \times k_{t}^{e}}$. It is called nested if for each $t \in T_{I} \backslash \mathcal{L}\left(T_{I}\right)$ there are transfer matrices $T_{e}^{t^{\prime} t} \in \mathbb{C}_{t^{\prime}}^{k^{\prime} \times k_{t}^{e}}$ such that for the restriction of the matrix $U_{e}(t)$ to the rows $t^{\prime}$ it holds that

$$
\begin{equation*}
\left.U_{e}(t)\right|_{t^{\prime}}=U_{e^{\prime}}\left(t^{\prime}\right) T_{e}^{t^{\prime} t} \quad \text { for all } t^{\prime} \in S_{I}(t) \tag{21}
\end{equation*}
$$

For estimating the complexity of storing a nested cluster basis $\mathcal{U}$, we assume that $k_{t}^{e} \leq k$ for all $t \in T_{I}, e \in \mathcal{E}(t)$ with some $k \in \mathbb{N}$. It follows from (13) that the depth $L \in \mathbb{N}$ of $T_{I}$ is of the order $L \sim \log _{2}\left(\kappa^{2}\right)$. Since the set of leaf clusters $\mathcal{L}\left(T_{I}\right)$ constitutes a partition of $I$ and according to (19) for each cluster $t \in \mathcal{L}\left(T_{I}\right)$ at most $|\mathcal{E}(t)| \sim \kappa^{2} 2^{-\ell}$ matrices $U_{e}(t)$ each with at most $k|t|$ entries have to be stored, $k|I| \kappa^{2} 2^{-L} \sim k N$ units of storage are required for the leaf matrices $U_{e}(t), t \in \mathcal{L}\left(T_{I}\right)$.

Using $|\mathcal{E}(t)| \leq 2\left|\mathcal{E}\left(t^{\prime}\right)\right|$ for $t^{\prime} \in S_{I}(t)$, we conclude that

$$
\sum_{t \in T_{I}^{(\ell)}}|\mathcal{E}(t)|=\sum_{t \in T_{I}^{(\ell)}} \frac{|\mathcal{E}(t)|}{2}\left|S_{I}(t)\right| \leq \sum_{t^{\prime} \in T_{I}^{(\ell+1)}}\left|\mathcal{E}\left(t^{\prime}\right)\right|,
$$

and with $\left|\mathcal{L}\left(T_{I}\right)\right| \leq|I| / n_{\text {min }}$ we can estimate the storage required for the transfer matrices

$$
k^{2} \sum_{t \in T_{I}}|\mathcal{E}(t)| \leq k^{2} L \sum_{t \in \mathcal{L}\left(T_{I}\right)}|\mathcal{E}(t)| \lesssim k^{2} L \kappa^{2} 2^{-L}\left|\mathcal{L}\left(T_{I}\right)\right| \lesssim \frac{k^{2}}{n_{\min }} N \log N .
$$

Definition 2. A matrix $M \in \mathbb{C}^{I \times J}$ is a directional $\mathcal{H}^{2}$-matrix if $M_{b}$ is of low rank for all $b \in P_{\text {low }}$ and there are nested directional cluster bases $\mathcal{U}$ and $\mathcal{V}$ such that for $t \times s \in P_{\text {high }}$

$$
\begin{equation*}
M_{t s}=U_{e}(t) S(t, s) V_{-e}^{H}(s) \tag{22}
\end{equation*}
$$

with coupling matrices $S(t, s) \in \mathbb{C}^{k_{t}^{e} \times k_{s}^{-e}}$ and $e \in \mathcal{E}(t)$ such that $s \subset F_{e}(t)$.
The storage cost of the blocks in $P_{\text {nonadm }}$ and $P_{\text {low }}$ is bounded by the storage cost of a hierarchical matrix, which is known to be $\mathcal{O}\left(\max \left\{k, n_{\min }\right\} N \log N\right)$. We estimate the storage required for the coupling matrices. According to Lemma 5, the number of blocks $t \times s \in P_{\text {high }}$ is $\mathcal{O}\left(2^{-\ell} \kappa^{2}\right)$ for $t \in T_{I}^{(\ell)}$. Thus, the coupling matrices require at most

$$
\begin{aligned}
k^{2} \sum_{t \in T_{I}}\left|\left\{s \in T_{J}: t \times s \in P_{\text {high }}\right\}\right| & \lesssim k^{2} \kappa^{2} \sum_{\ell=0}^{L-1} \sum_{t \in T_{I}^{(\ell)}} 2^{-\ell} \lesssim k^{2} \kappa^{2} L \sum_{t \in \mathcal{L}\left(T_{I}\right)} 2^{-L} \\
& =k^{2} \kappa^{2} 2^{-L} L\left|\mathcal{L}\left(T_{I}\right)\right| \lesssim \frac{k^{2}}{n_{\min }} N \log N
\end{aligned}
$$

units of storage. We obtain that the overall cost for a directional $\mathcal{H}^{2}$-matrix is of the order $\mathcal{O}\left(\max \left\{k^{2} / n_{\min }, k, n_{\min }\right\} N \log N\right)$, which is $\mathcal{O}(k N \log N)$ if $n_{\min }=k$.

### 4.1 Matrix-vector multiplication

Let $M \in \mathbb{C}^{I \times J}$ be a directional $\mathcal{H}^{2}$-matrix. Since its structure is similar to an $\mathcal{H}^{2}$-matrix, the matrixvector multiplication $y:=y+M x$ of $M$ by a vector $x \in \mathbb{C}^{J}$ can be done via the usual three-phase algorithm (cf. [27]) which we modified to account for the directions $e$. The following algorithm is a consequence of the decomposition

$$
M x=\sum_{t \times s \in P_{\text {nonadm }}} M_{t s} x_{s}+\sum_{t \times s \in P_{\text {low }}} X_{t s} Y_{t s}^{H} x_{s}+\sum_{t \times s \in P_{\text {high }}} U_{e}(t) S(t, s) V_{-e}(s)^{H} x_{s}
$$

with $e \in \mathcal{E}(t)$ such that $s \subset F_{e}(t)$.

## 1. Forward transform

The auxiliary vectors $\hat{x}_{e}(s):=V_{e}(s)^{H} x_{s}, e \in \mathcal{E}(s)$, are computed for all $s \in T_{J}$. Exploiting the nestedness of the cluster bases $\mathcal{V}$ (with transfer matrices $\bar{T}_{e}^{s^{\prime} s}$ ), one has the following recursive relation

$$
\hat{x}_{e}(s)=V_{e}(s)^{H} x_{s}=\sum_{s^{\prime} \in S_{J}(s)}\left(\bar{T}_{e}^{s^{\prime} s}\right)^{H} V_{e^{\prime}}\left(s^{\prime}\right)^{H} x_{s^{\prime}}=\sum_{s^{\prime} \in S_{J}(s)}\left(\bar{T}_{e}^{s^{\prime} s}\right)^{H} \hat{x}_{e^{\prime}}\left(s^{\prime}\right), \quad e \in \mathcal{E}(s)
$$

which has to be applied starting from the leaf vectors $\hat{x}_{e}(s), e \in \mathcal{E}(s), s \in \mathcal{L}\left(T_{J}\right)$.
2. Far field interaction

The products $S(t, s) \hat{x}_{-e}(s)$ are computed and summed up over all blocks $t \times s \in P_{\text {high }}$ :

$$
\hat{y}_{e}(t):=\sum_{s: t \times s \in P_{\mathrm{high}}} S(t, s) \hat{x}_{-e}(s), \quad e \in \mathcal{E}(t), t \in T_{I}
$$

## 3. Backward transform

The vectors $\hat{y}_{e}(t)$ are transformed to the target vector $y$. The nestedness (21) of $\mathcal{U}$ yields a recursion for the computation of $y:=\sum_{t \in T_{I}} \sum_{e \in \mathcal{E}(t)} U_{e}(t) \hat{y}_{e}(t)$, which descends $T_{I}$ :
(a) Compute $\hat{y}_{e^{\prime}}\left(t^{\prime}\right):=\hat{y}_{e^{\prime}}\left(t^{\prime}\right)+T_{e}^{t^{\prime} t} \hat{y}_{e}(t)$ for all $e \in \mathcal{E}(t)$ and all $t^{\prime} \in S_{I}(t)$;
(b) Compute $y_{t}:=y_{t}+\sum_{e \in \mathcal{E}(t)} U_{e}(t) \hat{y}_{e}(t)$ for all clusters $t \in \mathcal{L}\left(T_{I}\right)$.
4. Low-frequency interaction

For all $t \times s \in P_{\text {low }}$ compute $y_{t}:=y_{t}+X_{t s} Y_{t s}^{H} x_{s}$.
5. Near field interaction

For all $t \times s \in P_{\text {nonadm }}$ compute $y_{t}:=y_{t}+M_{t s} x_{s}$.
The total number of operations of the latter algorithm is bounded by $\mathcal{O}(k N \log N)$ for the choice $n_{\min }=k$, which can be proven with the same arguments as used for analyzing the storage complexity.

## 5 Construction of approximations

Our aim is to approximate $A \in \mathbb{C}^{I \times J}$ defined in (2) with a directional $\mathcal{H}^{2}$-matrix. Blocks in $P_{\text {nonadm }}$ are stored entrywise without approximation. From (7) it follows that $K$ is asymptotically smooth with constants independent of $\kappa$ on domains $X_{t} \times Y_{s}$ corresponding to blocks $t \times s \in P_{\text {low }}$. It follows from the convergence analysis in [7] that the adaptive cross approximation

$$
\begin{equation*}
A_{t s} \approx A_{t \sigma}\left(A_{\tau \sigma}\right)^{-1} A_{\tau s} \tag{23}
\end{equation*}
$$

with appropriately chosen $\tau \subset t$ and $\sigma \subset s, k_{\varepsilon}:=|\tau|=|\sigma| \sim|\log \varepsilon|^{2}$ independent of $\kappa$, can be used to generate low-rank approximations $X_{t s} Y_{t s}^{H} \approx A_{t s}$, where $X_{t s} \in \mathbb{C}^{t \times k_{\varepsilon}}$ and $Y_{t s} \in \mathbb{C}^{s \times k_{\varepsilon}}$.

In the rest of this section, we will consider the remaining case $t \times s \in P_{\text {high }}$, i.e. we assume that (15) is valid. To be able to apply the results from Sect. 2 and prove existence of low-rank matrix approximations it is required to additionally partition the far field $\mathcal{F}\left(X_{t}\right)$ into subsets $\mathcal{F}_{e}\left(X_{t}\right)$ corresponding to directions $e \in \mathcal{E}(t)$. Although ACA generates approximations of high quality, the number of blocks is too large (see Lemma 5) to construct and store the approximations as in the lowfrequency case. To overcome this difficulty, we consider the approximation (see [8] for the application of this kind of approximation to Laplace-type problems)

$$
\begin{equation*}
A_{t s} \approx A_{t \sigma_{t}^{e}}\left(A_{\tau_{t} \sigma_{t}^{e}}\right)^{-1} A_{\tau_{t} \sigma_{s}}\left(A_{\tau_{s}^{-e} \sigma_{s}}\right)^{-1} A_{\tau_{s}^{-e} s}, \quad Y_{s} \subset \mathcal{F}_{e}\left(X_{t}\right) \tag{24}
\end{equation*}
$$

instead of (23), which is of type (22) with coupling matrices $S(t, s)=A_{\tau_{t} \sigma_{s}}$. The aim of this section is to prove error estimates for the special type of low-rank approximation

$$
\begin{equation*}
A_{t s} \approx U_{e}(t) S(t, s) V_{-e}(s)^{H} \tag{25}
\end{equation*}
$$

with nested bases $\mathcal{U}$ and $\mathcal{V}$ approximating $A_{t \sigma_{t}^{e}}\left(A_{\tau_{t} \sigma_{t}^{e}}\right)^{-1}$ and $\left(A_{\tau_{s}^{-e} \sigma_{s}}\right)^{-1} A_{\tau_{s}^{-e} s}$, respectively.
A crucial part of the approximation (24) is the construction of what we call proper pivots $\tau_{t} \subset t$ and $\sigma_{t}^{e} \subset F_{e}(t),\left|\tau_{t}\right|=\left|\sigma_{t}^{e}\right|$. They have to guarantee that $A_{\tau_{t} \sigma_{t}^{e}}$ is invertible, and for proving error estimates they have to be chosen so that

$$
\left\|A_{t s}-A_{t \sigma_{t}^{e}}\left(A_{\tau_{t} \sigma_{t}^{e}}\right)^{-1} A_{\tau_{t} s}\right\| \leq c_{R} \varepsilon\left\|A_{t J}\right\| \quad \text { for all } s \subset F_{e}(t)
$$

with some $c_{R}>0$; cf. Lemma 9. Hence $\tau_{t}$ and $\sigma_{t}^{e}$ represent $t$ and its far field $F_{e}(t)$, respectively. We refer to [8] for a method for choosing $\tau_{t}$ and $\sigma_{t}^{e}$. Note that it is sufficient to choose $\sigma_{t}^{e}$ so that

$$
\begin{equation*}
Y_{\sigma_{t}^{e}} \subset\left(\chi\left(X_{t}\right)+Z_{e}\right) \cap \mathcal{F}\left(X_{t}\right) \tag{26}
\end{equation*}
$$

In the sequel, $\varepsilon>0$ denotes a given accuracy that (up to constants) is to be achieved by the approximations. Let $\left\{\zeta_{1}, \ldots, \zeta_{k_{\varepsilon}}\right\}$ be a basis of the space

$$
\hat{\Pi}_{p_{\varepsilon}-1}^{3}:=\left\{u \exp (2 \pi \mathrm{i} \kappa(\cdot, e)): u \in \Pi_{p_{\varepsilon}-1}^{3}\right\}
$$

where $k_{\varepsilon}:=\operatorname{dim} \Pi_{p_{\varepsilon}-1}^{3} \sim p_{\varepsilon}^{3}$ and $p_{\varepsilon} \in \mathbb{N}$ is the smallest number such that $p_{\varepsilon} \geq\left|\log _{\rho} \varepsilon\right|$ with $\rho$ from (11).

Assumption 1. Let $t \in T_{I}$. If $|t| \geq k_{\varepsilon}$, we assume that there is $\tau_{t}=\left\{i_{1}, \ldots, i_{k_{\varepsilon}}\right\} \subset t$ such that the following two conditions are satisfied.
(i) There are coefficients $\xi_{i \ell}$ such that

$$
\begin{equation*}
\left(\varphi_{i}, \zeta_{j}\right)_{L^{2}(\Gamma)}=\sum_{\ell=1}^{k_{\varepsilon}} \xi_{i \ell}\left(\varphi_{i_{\ell}}, \zeta_{j}\right)_{L^{2}(\Gamma)}, \quad i \in t, j=1, \ldots, k_{\varepsilon} \tag{27}
\end{equation*}
$$

and the norm of the matrix $\Xi=\left(\xi_{i \ell}\right)_{i \ell} \in \mathbb{C}^{t \times k_{\varepsilon}}$ is bounded by a multiple of $2^{p_{\varepsilon}}$ provided that $\left\|\varphi_{i}\right\|_{L^{1}(\Gamma)}=1, i \in t$,
(ii) each matrix $A_{\tau_{t} F_{e}(t)}, e \in \mathcal{E}(t)$, has full rank.

In the remaining case $|t|<k_{\varepsilon}$, we set $\tau_{t}=t$ and assume that $A_{t F_{e}(t)}$ has full rank.
Notice that with the previous assumptions it is possible to guarantee that the rows $\tau_{t}$ used for the approximation of $A_{t s}$ can be chosen independently of $s \subset F_{e}(t)$. This will be crucial for the construction of nested bases.

In the following lemma, we prove error estimates for the multivariate tensor product Chebyshev interpolant $\Im_{p} \hat{K}(\cdot, y) \in \Pi_{p}^{3}$ of $\hat{K}(\cdot, y):=K(\cdot, y) \exp (-2 \pi \mathrm{i} \kappa(\cdot-y, e))$ with fixed $y \in Y$.

Lemma 7. Let $X, Y \subset \mathbb{R}^{3}$ such that $Y \subset \mathcal{F}_{e}(X)$. Then

$$
\left|\hat{K}(x, y)-\mathfrak{I}_{x, p-1} \hat{K}(x, y)\right| \leq c_{\mathfrak{I}}(p)\left(\frac{\rho}{\rho+2}\right)^{p} \max _{x^{\prime} \in X}\left|\hat{K}\left(x^{\prime}, y\right)\right| \quad \text { for all } x \in X, y \in Y
$$

where $c_{\mathfrak{I}}(p):=8 \operatorname{ecp}(1+\rho)\left(1+\frac{2}{\pi} \log p\right)^{3}$.
Proof. Without loss of generality we may assume that $X$ is contained in a cube $Q=\prod_{i=1}^{3} Q_{i}$ which is aligned with $e$ and that $Y \subset \mathcal{F}_{e}(Q)$. Notice that this can be achieved by slightly modifying the constants $\gamma_{\text {high }}, \eta_{\text {high }}$. Let $\hat{K}_{i}$ be the function in the $i$-th argument of $\hat{K}(\cdot, y), i=1,2,3$. From (11) we obtain

$$
\left\|\hat{K}_{i}^{(k)}\right\|_{Q_{i}, \infty} \leq c k!\left(\frac{\rho}{\operatorname{diam} Q_{i}}\right)^{k}\left\|\hat{K}_{i}\right\|_{Q_{i}, \infty}, \quad k \in \mathbb{N}
$$

Using [9, Lemma 3.13], this implies

$$
\min _{q \in \Pi_{p-1}}\left\|\hat{K}_{i}-q\right\|_{Q_{i}, \infty} \leq \tilde{c} p\left(\frac{\rho}{\rho+2}\right)^{p}\left\|\hat{K}_{i}\right\|_{Q_{i}, \infty}
$$

where $\tilde{c}:=4 \mathrm{e} c(1+\rho)$. With this estimate the proof can be done analogously to Theorem 3.18 in [7].

Although most of the estimates will hold also in other matrix norms, throughout this article the maximum absolute column sum

$$
\|A\|:=\max _{j=1, \ldots, n} \sum_{i=1}^{m}\left|a_{i j}\right|
$$

of $A \in \mathbb{C}^{m \times n}$ will be used if not otherwise indicated.
Lemma 8. Let assumption (i) be valid and let $\varphi_{i}, \psi_{j}$, and $f$ in (2) be non-negative. Assume that $c_{\mathrm{as}, 2} \eta_{\text {high }}<1$. For $t \in T_{I}$ satisfying $|t| \geq k_{\varepsilon}$ and $e \in \mathcal{E}(t)$ there is $\Xi \in \mathbb{R}^{t \times k_{\varepsilon}}$ and $c_{1}>0$ such that

$$
\left\|A_{t s}-\Xi A_{\tau_{t} s}\right\| \leq c_{1} \varepsilon\left\|A_{t s}\right\|
$$

for all $s \subset J$ satisfying $Y_{s} \subset \mathcal{F}_{e}\left(X_{t}\right)$.
Proof. Due to $K(x, y)=\exp (2 \pi \mathrm{i} \kappa(x-y, e)) \hat{K}(x, y)$, we can apply Lemma 7 and obtain for $x \in X_{t}$, $y \in Y_{s}$ that with $T_{p_{\varepsilon}}(x, y):=\exp (2 \pi \mathrm{i} \kappa(x, e)) \Im_{x, p_{\varepsilon}-1} \hat{K}(x, y)$

$$
\left|K(x, y)-T_{p_{\varepsilon}}(x, y)\right|=\left|\hat{K}(x, y)-\Im_{x, p_{\varepsilon}-1} \hat{K}(x, y)\right| \leq c_{\mathfrak{J}}\left(p_{\varepsilon}\right)\left(\frac{\rho}{\rho+2}\right)^{p_{\varepsilon}} \max _{x^{\prime} \in X_{t}}\left|\hat{K}\left(x^{\prime}, y\right)\right|
$$

Without loss of generality we may assume that $\left\|\varphi_{i}\right\|_{L^{1}(\Gamma)}=1$. Then assumption (27) is equivalent with

$$
\int_{\Gamma}\left(\varphi_{i}(x)-\sum_{\ell=1}^{k_{\varepsilon}} \xi_{i \ell} \varphi_{i_{\ell}}(x)\right) \zeta(x) \mathrm{d} s_{x}=0 \quad \text { for all } \zeta \in \hat{\Pi}_{p_{\varepsilon}-1}^{3}
$$

Defining the matrix $\Xi \in \mathbb{R}^{t \times k_{\varepsilon}}$ with the entries $\xi_{i \ell}$, from

$$
\begin{aligned}
a_{i j}-\sum_{\ell=1}^{k_{\varepsilon}} \xi_{i \ell} a_{i_{\ell} j} & =\int_{\Gamma} \int_{\Gamma}\left(\varphi_{i}(x)-\sum_{\ell=1}^{k_{\varepsilon}} \xi_{i \ell} \varphi_{i_{\ell}}(x)\right) K(x, y) \psi_{j}(y) \mathrm{d} s_{y} \mathrm{~d} s_{x} \\
& =\int_{\Gamma} \int_{\Gamma}\left(\varphi_{i}(x)-\sum_{\ell=1}^{k_{\varepsilon}} \xi_{i \ell} \varphi_{i_{\ell}}(x)\right)\left[K(x, y)-T_{p_{\varepsilon}}(x, y)\right] \psi_{j}(y) \mathrm{d} s_{y} \mathrm{~d} s_{x}
\end{aligned}
$$

we see that

$$
\left|a_{i j}-\sum_{\ell=1}^{k_{\varepsilon}} \xi_{i \ell} a_{i_{\ell} j}\right| \leq \tilde{c}\left(\frac{\rho}{\rho+2}\right)^{p_{\varepsilon}} \max _{x^{\prime} \in X_{t}} \int_{\Gamma}\left|\hat{K}\left(x^{\prime}, y\right)\right|\left|\psi_{j}(y)\right| \mathrm{d} s_{y}
$$

with $\tilde{c}:=c_{\mathfrak{I}}\left(p_{\varepsilon}\right)\left(1+\|\Xi\|_{\infty}\right)$. From the Taylor expansion and the asymptotic smoothness of $f$ it can be seen that for $c_{\mathrm{as}, 2} \eta_{\text {high }}<1$

$$
\begin{equation*}
\left|f\left(x^{\prime}, y\right)\right| \leq \hat{c}|f(x, y)| \quad \text { for all } x, x^{\prime} \in X_{t} \tag{28}
\end{equation*}
$$

with a constant $\hat{c}>0$. Estimate (28) and $\left\|\varphi_{i}\right\|_{L^{1}(\Gamma)}=1$ imply for $y \in Y_{s}$

$$
\max _{x^{\prime} \in X_{t}}\left|\hat{K}\left(x^{\prime}, y\right)\right|=\max _{x^{\prime} \in X_{t}}\left|f\left(x^{\prime}, y\right)\right| \leq \hat{c} \min _{x \in X_{i}}|f(x, y)| \leq \hat{c} \int_{\Gamma}\left|\varphi_{i}(x)\right||f(x, y)| \mathrm{d} s_{x}
$$

From

$$
\begin{aligned}
\left|a_{i j}\right| & =\left|\mathrm{e}^{2 \pi \mathrm{i} \kappa\left[-\left|\xi\left(X_{i}\right)-\xi\left(Y_{j}\right)\right|\right]} a_{i j}\right| \geq \operatorname{Re} \int_{\Gamma} \int_{\Gamma} \varphi_{i}(x) \psi_{j}(y) f(x, y) \mathrm{e}^{2 \pi \mathrm{i} \kappa\left[|x-y|-\left|\xi\left(X_{i}\right)-\xi\left(Y_{j}\right)\right|\right]} \mathrm{d} s_{y} \mathrm{~d} s_{x} \\
& =\int_{\Gamma} \int_{\Gamma} \varphi_{i}(x) \psi_{j}(y) f(x, y) \cos \left(2 \pi \kappa\left[|x-y|-\left|\xi\left(X_{i}\right)-\xi\left(Y_{j}\right)\right|\right]\right) \mathrm{d} s_{y} \mathrm{~d} s_{x}
\end{aligned}
$$

$$
\geq c^{\prime} \int_{\Gamma} \int_{\Gamma} \varphi_{i}(x) \psi_{j}(y) f(x, y) \mathrm{d} s_{y} \mathrm{~d} s_{x}
$$

with $c^{\prime}>0$ independent of $\kappa$ and $h$ due to

$$
\kappa\left[|x-y|-\left|\xi\left(X_{i}\right)-\xi\left(Y_{j}\right)\right|\right] \leq \kappa\left[\left|x-\xi\left(X_{i}\right)\right|+\left|y-\xi\left(Y_{j}\right)\right|\right] \leq 2 \kappa h \leq 2 c_{\kappa}<\frac{1}{2}
$$

we obtain exploiting the non-negativity of $\varphi_{i}, \psi_{j}$, and $f$

$$
\left|a_{i j}-\sum_{\ell=1}^{k_{\varepsilon}} \xi_{i \ell} a_{i_{\ell} j}\right| \leq \frac{\hat{c} \tilde{c}}{c^{\prime}}\left(\frac{\rho}{\rho+2}\right)^{p_{\varepsilon}}\left|a_{i j}\right|
$$

Hence, the matrix $\Xi$ satisfies

$$
\left\|A_{t s}-\Xi A_{\tau_{t} s}\right\| \leq \frac{\hat{c} \tilde{c}}{c^{\prime}(\rho+2)^{p_{\varepsilon}}} \rho^{p_{\varepsilon}}\left\|A_{t s}\right\| \leq c_{1} \varepsilon\left\|A_{t s}\right\|
$$

because $\frac{\hat{c} c_{\mathcal{J}}\left(p_{\varepsilon}\right)}{c^{\prime}(\rho+2)^{p_{\varepsilon}}}\left(1+\|\Xi\|_{\infty}\right)$ is bounded by a constant $c_{1}$ from above due to the assumption that $\|\Xi\|_{\infty}$ is bounded by a multiple of $2^{p_{\varepsilon}}$.

The expression

$$
c_{S}:=\max \left\{\left\|\left(A_{\tau_{t} \sigma_{t}^{e}}\right)^{-1} A_{\tau_{t} s}\right\|: s \subset J, Y_{s} \subset \mathcal{F}_{e}\left(X_{t}\right), e \in \mathcal{E}(t), t \in T_{I}\right\}
$$

will play a central role in the following error analysis. Note that $c_{S}$ depends on the choice of the pivots $\tau_{t}, \sigma_{t}^{e}$.

Lemma 9. Let Assumption 1 be valid. Then for $t \in T_{I}$ and $e \in \mathcal{E}(t)$ there are proper pivots $\left(\tau_{t}, \sigma_{t}^{e}\right)$, $\left|\tau_{t}\right|=\left|\sigma_{t}^{e}\right|=\min \left\{k_{\varepsilon},|t|\right\}$, i.e., for all $s \subset J$ satisfying $Y_{s} \subset \mathcal{F}_{e}\left(X_{t}\right)$

$$
\begin{equation*}
\left\|A_{t s}-A_{t \sigma_{t}^{e}}\left(A_{\tau_{t} \sigma_{t}^{e}}\right)^{-1} A_{\tau_{t} s}\right\| \leq c_{2} \varepsilon\left\|A_{t J}\right\| \tag{29}
\end{equation*}
$$

where $c_{2}:=c_{1}\left(1+c_{S}\right)$.
Proof. Since $A_{t F_{e}(t)}$ has full rank, there is $\sigma_{t}^{e} \subset F_{e}(t),\left|\sigma_{t}^{e}\right|=\left|\tau_{t}\right|=\min \left\{k_{\varepsilon},|t|\right\}$, such that $A_{\tau_{t} \sigma_{t}^{e}}$ is invertible. If $|t|<k_{\varepsilon}$, we have $\tau_{t}=t$ and $A_{t \sigma_{t}^{e}}\left(A_{\tau_{t} \sigma_{t}^{e}}\right)^{-1}=I$. Hence, (29) is trivially satisfied. We may hence assume that $|t| \geq k_{\varepsilon}$. Let $\Xi \in \mathbb{R}^{t \times k_{\varepsilon}}$ be as in Lemma 8. We have

$$
A_{t s}-A_{t \sigma_{t}^{e}}\left(A_{\tau_{t} \sigma_{t}^{e}}\right)^{-1} A_{\tau_{t} s}=\left\{A_{t s}-\Xi A_{\tau_{t} s}\right\}-\left\{A_{t \sigma_{t}^{e}}-\Xi A_{\tau_{t} \sigma_{t}^{e}}\right\}\left(A_{\tau_{t} \sigma_{t}^{e}}\right)^{-1} A_{\tau_{t} s}
$$

and thus

$$
\begin{aligned}
\left\|A_{t s}-A_{t \sigma_{t}^{e}}\left(A_{\tau_{t} \sigma_{t}^{e}}\right)^{-1} A_{\tau_{t} s}\right\| & \leq\left\|A_{t s}-\Xi A_{\tau_{t} s}\right\|+c_{S}\left\|A_{t \sigma_{t}^{e}}-\Xi A_{\tau_{t} \sigma_{t}^{e}}\right\| \\
& \leq c_{1} \varepsilon\left(\left\|A_{t s}\right\|+c_{S}\left\|A_{t \sigma_{t}^{e}}\right\|\right) \leq c_{2} \varepsilon\left\|A_{t J}\right\|
\end{aligned}
$$

The second last estimate follows from Lemma 8 , because $Y_{\sigma_{t}^{e}}, Y_{s} \subset \mathcal{F}_{e}\left(X_{t}\right)$.

### 5.1 Construction of directional cluster bases

Based on the previous estimates, we are going to construct and analyze nested bases approximations. The construction of nested bases is usually done by explicit approximation of the kernel function $K$; see, for instance, the fast multipole method [32] and the method in [29], which uses interpolation. In this section, we construct the nested bases via a purely algebraic technique which is based on the original matrix entries and thus avoids explicit kernel expansions. In this sense, the presented construction is in the class of kernel independent fast multipole methods; see [3, 34, 17].
We will define a nested basis $\mathcal{U}$ consisting of matrices $U_{e}(t) \in \mathbb{C}^{t \times k_{\varepsilon}}$ for each $t \in T_{I}$ and $e \in \mathcal{E}(t)$ in a recursive manner starting from the leaves. Due to (15a), it is actually sufficient to consider the sub-tree

$$
\hat{T}_{I}:=\left\{t \in T_{I}: \kappa \operatorname{diam} X_{t}>1\right\}
$$

of $T_{I}$. For its leaf clusters $t \in \mathcal{L}\left(\hat{T}_{I}\right)$ and $e \in \mathcal{E}(t)$ we set $U_{e}(t)=B_{e}^{t t}$, where for $t^{\prime} \subset t$

$$
B_{e}^{t^{\prime} t}:=A_{t^{\prime} \sigma_{t}^{e}}\left(A_{\tau_{t} \sigma_{t}^{e}}\right)^{-1} .
$$

Assume that matrices $U_{e}\left(t^{\prime}\right)$ have already been constructed for the sons $t^{\prime} \in S_{I}(t)$ of $t \in \hat{T}_{I} \backslash \mathcal{L}\left(\hat{T}_{I}\right)$ and $e \in \mathcal{E}\left(t^{\prime}\right)$. Then in view of (21) we define for $e \in \mathcal{E}(t)$

$$
\left.U_{e}(t)\right|_{t^{\prime}}:=U_{e^{\prime}}\left(t^{\prime}\right) B_{e^{\tau^{\prime}} t}, \quad t^{\prime} \in S_{I}(t),
$$

where $e^{\prime} \in \mathcal{E}\left(t^{\prime}\right)$ is defined before Lemma 6 .
The following lemma estimates the accuracy when expressing $B_{e}^{t^{\prime} t}$ by the product $B_{e^{\prime}}^{t^{\prime} t^{\prime}} B_{e}^{\tau_{t^{\prime}} t}$. As stated in [8], $B_{e}^{t^{\prime} t}$ may be regarded as the algebraic form of an interpolation operator.

Lemma 10. Let $t^{\prime} \in \hat{T}_{I}$ satisfy $t^{\prime} \subset t \in \hat{T}_{I} \backslash \mathcal{L}\left(\hat{T}_{I}\right)$ and $e \in \mathcal{E}(t)$. Then for all $s \subset J$ satisfying $Y_{s} \subset \mathcal{F}_{e}\left(X_{t}\right)$ it holds that

$$
\left\|\left[B_{e}^{t^{\prime} t}-B_{e^{\prime}}^{t^{\prime} t^{\prime}} B_{e}^{\tau_{t^{\prime}} t}\right] A_{\tau_{t} s}\right\| \leq c_{2} c_{S} \varepsilon\left\|A_{t^{\prime} J}\right\| .
$$

Proof. It is clear that

$$
\left[B_{e}^{t^{\prime} t}-B_{e^{\prime}}^{t^{\prime} t^{\prime}} B_{e}^{\tau_{t} t}\right] A_{\tau_{t} s}=\left[A_{t^{\prime} \sigma_{t}^{e}}-A_{t^{\prime} \sigma_{t^{\prime}}^{e^{\prime}}}\left(A_{\tau_{t^{\prime}} \sigma_{t^{\prime}}^{\sigma^{\prime}}}\right)^{-1} A_{\tau_{t^{\prime}} \sigma_{t}^{e}}\right]\left(A_{\tau_{t} \sigma_{t}^{e}}\right)^{-1} A_{\tau_{t} s}
$$

Due to (26) and (20), it holds that $Y_{\sigma_{t}^{e}} \subset \mathcal{F}_{e^{\prime}}\left(X_{t^{\prime}}\right)$. Hence, Lemma 9 yields

Theorem 1. Let Assumption 1 be valid. Let $t \in \hat{T}_{I}$ and let $\ell=L\left(\hat{T}_{t}\right)$ denote the depth of the sub-tree $\hat{T}_{t}$ of $\hat{T}_{I}$ rooted at $t \in T_{I}$. Then for all $e \in \mathcal{E}(t)$

$$
\left\|\left[U_{e}(t)-B_{e}^{t t}\right] A_{\tau_{t} s}\right\| \leq c_{3} \varepsilon\left\|A_{t J}\right\| \quad \text { for all } s \subset F_{e}(t)
$$

where

$$
c_{3}:=c_{1}\left(1+c_{S}\right) \frac{\left(2 c_{S}\right)^{\ell}}{2 c_{S}-1} .
$$

Proof. From Lemma 10 we have for $t \in \hat{T}_{I} \backslash \mathcal{L}\left(\hat{T}_{I}\right)$

$$
\begin{aligned}
\|\left[U_{e}(t)\right. & \left.-B_{e}^{t t}\right] A_{\tau_{t} s}\left\|\leq \sum_{t^{\prime} \in S_{I}(t)}\right\|\left[U_{e}(t) \mid t^{\prime}-B_{e}^{t^{\prime} t}\right] A_{\tau_{t} s}\left\|=\sum_{t^{\prime} \in S_{I}(t)}\right\|\left[U_{e^{\prime}}\left(t^{\prime}\right) B_{e}^{\tau_{t^{\prime}} t}-B_{e}^{t^{\prime} t}\right] A_{\tau_{t} s} \| \\
& \leq \sum_{t^{\prime} \in S_{I}(t)}\left\|\left[U_{e^{\prime}}\left(t^{\prime}\right)-B_{e^{\prime} t^{\prime}}^{t^{\prime}}\right] B_{e}^{\tau_{t^{\prime} t} t} A_{\tau_{t s} s}\right\|+\left\|\left[B_{e}^{t^{\prime} t}-B_{e^{\prime}}^{t^{\prime} t^{\prime}} B_{e}^{\tau_{t^{\prime}} t}\right] A_{\tau_{t} s}\right\| \\
& \leq \sum_{t^{\prime} \in S_{I}(t)} c_{S}\left\|\left[U_{e^{\prime}}\left(t^{\prime}\right)-B_{e^{\prime}}^{t^{\prime} t^{\prime}}\right] A_{\tau_{t^{\prime}} \sigma_{t}^{e}}\right\|+\varepsilon c_{2} c_{S}\left\|A_{t^{\prime} J}\right\| \\
& \leq 2 c_{2} c_{S} \varepsilon\left\|A_{t J}\right\|+c_{S} \sum_{t^{\prime} \in S_{I}(t)}\left\|\left[U_{e^{\prime}}\left(t^{\prime}\right)-B_{e^{\prime}}^{t^{\prime} t^{\prime}}\right] A_{\tau_{t^{\prime}} \sigma_{t}^{e}}\right\| .
\end{aligned}
$$

We set $\alpha_{t^{\prime}}:=\left\|\left[U_{e^{\prime}}\left(t^{\prime}\right)-B_{e^{\prime}}^{t^{\prime} t^{\prime}}\right] A_{\tau_{t^{\prime}} \sigma_{t}^{e}}\right\|$ for $t^{\prime} \in S_{I}(t)$. From (26) and (20) we obtain that $Y_{\sigma_{t}^{e}} \subset \mathcal{F}_{e^{\prime}}\left(X_{t^{\prime}}\right)$. Hence, we can apply the latter inequality recursively replacing $s$ by $\sigma_{t}^{e}$ and obtain the recurrence relation

$$
\begin{equation*}
\alpha_{t} \leq 2 c_{2} c_{S} \varepsilon\left\|A_{t J}\right\|+c_{S} \sum_{t^{\prime} \in S_{I}(t)} \alpha_{t^{\prime}}, \quad t \in T_{I} \backslash \mathcal{L}\left(T_{I}\right) . \tag{30}
\end{equation*}
$$

We show that

$$
\begin{equation*}
\alpha_{t} \leq 2 c_{2} c_{S} \varepsilon \frac{\left(2 c_{S}\right)^{\ell-1}-1}{2 c_{S}-1}\left\|A_{t J}\right\|, \quad t \in T_{I} \tag{31}
\end{equation*}
$$

where $\ell=\ell(t)$ denotes the depth of the sub-tree $\hat{T}_{t}$. This estimate is obviously true for leaf clusters $t \in \mathcal{L}\left(T_{I}\right)$ as $\alpha_{t}=0$. Assume that (31) is valid for the sons $S_{I}(t)$ of $t \in \hat{T}_{I} \backslash \mathcal{L}\left(\hat{T}_{I}\right)$. Then (30) proves

$$
\begin{aligned}
\alpha_{t} & \leq 2 c_{2} c_{S} \varepsilon\left\|A_{t J}\right\|+c_{S} \sum_{t^{\prime} \in S_{I}(t)} \alpha_{t^{\prime}} \leq 2 c_{2} c_{S} \varepsilon\left\|A_{t J}\right\|+2 c_{2} c_{S}^{2} \varepsilon \frac{\left(2 c_{S}\right)^{\ell-2}-1}{2 c_{S}-1} \sum_{t^{\prime} \in S_{I}(t)}\left\|A_{t^{\prime} J}\right\| \\
& \leq 2 c_{2} c_{S} \varepsilon\left(1+2 c_{S} \frac{\left(2 c_{S}\right)^{\ell-2}-1}{2 c_{S}-1}\right)\left\|A_{t J}\right\|=2 c_{2} c_{S} \varepsilon \frac{\left(2 c_{S}\right)^{\ell-1}-1}{2 c_{S}-1}\left\|A_{t J}\right\| .
\end{aligned}
$$

Hence,

$$
\begin{aligned}
\left\|\left[U_{e}(t)-B_{e}^{t t}\right] A_{\tau_{t} s}\right\| & \leq 2 c_{2} c_{S} \varepsilon\left\|A_{t J}\right\|+c_{S} \sum_{t^{\prime} \in S_{I}(t)} \alpha_{t^{\prime}} \\
& \leq 2 c_{2} c_{S} \varepsilon \frac{\left(2 c_{S}\right)^{\ell-1}}{2 c_{S}-1}\left\|A_{t J}\right\| \leq c_{2} \varepsilon \frac{\left(2 c_{S}\right)^{\ell}}{2 c_{S}-1}\left\|A_{t J}\right\|
\end{aligned}
$$

together with $c_{2}=c_{1}\left(1+c_{S}\right)$ proves the assertion.
Similar results as for the row clusters $t$ can be obtained for column clusters $s \in T_{J}$ and $e \in \mathcal{E}(s)$ provided assumptions analogous to Assumption 1 are made. In particular, this defines clusters $\sigma_{s} \subset s$ and $\tau_{s}^{e} \subset F_{e}^{\prime}(s),\left|\tau_{s}^{e}\right|=\left|\sigma_{s}\right|$, where

$$
F_{e}^{\prime}(s):=\bigcup\left\{t \in T_{I}: \exists \hat{s} \supset s \text { such that } t \times \hat{s} \in P_{\text {high }}\right\} \cap\left\{i \in I: X_{i} \subset \mathcal{F}_{e}\left(Y_{s}\right)\right\}
$$

For $s^{\prime} \subset s$ we define the begin

$$
C_{e}^{s^{\prime} s}:=\left(A_{\tau_{s}^{e} \sigma_{s}}\right)^{-1} A_{\tau_{s}^{e} s^{\prime}} .
$$

For leaf clusters $s \in \mathcal{L}\left(\hat{T}_{J}\right)$ and $e \in \mathcal{E}(s)$ we set $V_{e}(s)=\left(C_{e}^{s s}\right)^{H}$. Assuming that matrices $V_{e}\left(s^{\prime}\right)$ have already been constructed for the sons $s^{\prime} \in S_{J}(s)$ of $s \in \hat{T}_{J} \backslash \mathcal{L}\left(\hat{T}_{J}\right)$ and $e \in \mathcal{E}\left(s^{\prime}\right)$, we define for $e \in \mathcal{E}(s)$

$$
\left.V_{e}(t)\right|_{s^{\prime}}:=V_{e^{\prime}}\left(s^{\prime}\right)\left(C_{e}^{\tau_{\tau^{\prime}} t}\right)^{T}, \quad s^{\prime} \in S_{J}(s),
$$

where $e^{\prime} \in \mathcal{E}\left(s^{\prime}\right)$ is defined before Lemma 6 . Due to the analogy, we omit the proofs of the following error estimates.

Lemma 11. Then for $s \in T_{J}$ and $e \in \mathcal{E}(s)$ there are proper pivots $\left(\tau_{s}^{e}, \sigma_{t}\right),\left|\tau_{t}^{e}\right|=\left|\sigma_{t}\right|=\min \left\{k_{\varepsilon},|s|\right\}$, i.e.,

$$
\left\|A_{t s}-A_{t \sigma_{s}}\left(A_{\tau_{s}^{e} \sigma_{s}}\right)^{-1} A_{\tau_{s}^{e} s}\right\| \leq c_{4} \varepsilon\left\|A_{I s}\right\|
$$

for all $t \subset I$ satisfying $X_{t} \subset \mathcal{F}_{e}\left(Y_{s}\right)$.
Theorem 2. Let $s \in \hat{T}_{J}$ and let $\ell=L\left(\hat{T}_{s}\right)$ denote the depth of the cluster tree $\hat{T}_{s}$. Then there is $c_{5}>0$ such that for all $e \in \mathcal{E}(s)$

$$
\left\|A_{t \sigma_{s}}\left[C_{e}^{s s}-V_{e}(s)^{H}\right]\right\| \leq c_{5} \varepsilon\left\|A_{I s}\right\|
$$

for all $t \subset F_{e}^{\prime}(s)$.
Using the previously constructed bases $\mathcal{U}$ and $\mathcal{V}$, we employ $S(t, s):=A_{\tau_{t} \sigma_{s}}$ in (25) for each block $A_{t s}, t \times s \in P_{\text {high }}$. In the following theorem, the accuracy of the nested approximation based on the matrix entries of $A$ is estimated.

Theorem 3. Let all previous assumptions be valid and $t \times s \in P_{\text {high }}$. Then there is $e \in \mathcal{E}(t)$ such that $s \subset F_{e}(t)$ and the approximation error is bounded by

$$
\left\|A_{t s}-U_{e}(t) S(t, s) V_{-e}(s)^{H}\right\| \leq\left[c_{2}+c_{3}\left\|C_{-e}^{s s}\right\|\right] \varepsilon\left\|A_{t J}\right\|+\left[c_{4}\left\|B_{e}^{t t}\right\|+c_{5}\left\|U_{e}(t)\right\|\right] \varepsilon\left\|A_{I s}\right\| .
$$

Proof. From $s \subset F_{e}(t)$ it follows that $t \subset F_{-e}^{\prime}(s)$. We have that

$$
A_{t s}-B_{e}^{t t} S(t, s) C_{-e}^{s s}=A_{t s}-B_{e}^{t t} A_{\tau_{t} s}+B_{e}^{t t}\left[A_{\tau_{t} s}-A_{\tau_{t} \sigma_{s}} C_{-e}^{s s}\right] .
$$

From Lemma 9 it follows that $\left\|A_{t s}-B_{e}^{\text {tt }} A_{\tau_{t} s}\right\| \leq c_{2} \varepsilon\left\|A_{t J}\right\|$, and from Lemma 11 we have that $\left\|A_{\tau_{t} s}-A_{\tau_{t} \sigma_{s}} C_{-e}^{s s}\right\| \leq c_{4} \varepsilon\left\|A_{I s}\right\|$. Therefore,

$$
\left\|A_{t s}-B_{e}^{t t} S(t, s) C_{-e}^{s s}\right\| \leq \varepsilon\left[c_{2}\left\|A_{t J}\right\|+c_{4}\left\|B_{e}^{t t}\right\|\left\|A_{I s}\right\|\right] .
$$

Furthermore, Theorems 1 and 2 yield

$$
\begin{aligned}
\| U_{e}(t) S(t, s) V_{-e}(s)^{H} & -B_{e}^{t t} S(t, s) C_{-e}^{s s} \| \\
& \leq\left\|U_{e}(t)\right\|\left\|S(t, s)\left[V_{e}(s)^{H}-C_{-}^{s s}\right]\right\|+\left\|\left[U_{e}(t)-B_{e}^{t t}\right] S(t, s)\right\|\left\|C_{-e}^{s s}\right\| \\
& \leq c_{5} \varepsilon\left\|U_{e}(t)\right\|\left\|A_{I s}\right\|+c_{3} \varepsilon\left\|C_{-e}^{s s}\right\|\left\|A_{t J}\right\|,
\end{aligned}
$$

which proves the assertion.

## 6 Numerical Results

We consider the sound soft/hard scattering problem (1) imposing either the Dirichlet datum $u=u_{D}$ (sound soft) or the Neumann trace $\partial_{\nu} u=v_{N}$ (sound hard) on the boundary $\Gamma$. We denote $\mathfrak{V}$ the single and $\mathfrak{K}$ the double-layer-operator with the kernels $S(x-y)$ and $\partial_{\nu_{y}} S(x-y)$, respectively. Using Green's representation formula $u=\mathfrak{K} u-\mathfrak{V}\left(\partial_{\nu} u\right)$ in $\Omega^{c}$ and the jump relations, we end up with the integral equation

$$
\begin{equation*}
\mathfrak{V}\left(\partial_{\nu} u\right)=\left(\mathfrak{K}-\frac{1}{2} \mathfrak{I}\right) u \quad \text { on } \Gamma, \tag{32}
\end{equation*}
$$

which can be solved either for the unknown $\left.\partial_{\nu} u\right|_{\Gamma}$ in the case of a Dirichlet problem or for the unknown $\left.u\right|_{\Gamma}$ in the Neumann case. The Brakhage-Werner ansatz

$$
\begin{equation*}
u=\mathfrak{K} \phi-\mathrm{i} \beta \mathfrak{V} \phi \quad \text { in } \Omega^{c} \tag{33}
\end{equation*}
$$

with an arbitrary coupling parameter $\beta>0$ uses an unknown density function $\phi$ that satisfies the integral equation

$$
\begin{equation*}
\left(\frac{1}{2} \mathfrak{I}+\mathfrak{K}-\mathrm{i} \beta \mathfrak{V}\right) \phi=u_{D} \tag{34}
\end{equation*}
$$

In either case, Galerkin discretization of the integral equations leads to a linear system with matrices of the form (2) that can be approximated by directional $\mathcal{H}^{2}$-matrices. The solution can then be obtained via GMRES using the matrix-vector multiplication from Sect. 4.1, which we proved to have logarithmic-linear complexity.

## Approximation of $\mathfrak{V}$

As a first step, we validate the logarithmic-linear complexity of the directional $\mathcal{H}^{2}$-matrices (labeled $\left.\operatorname{dir} \mathcal{H}^{2}-\mathrm{ACA}\right)$. Moreover, we compare our new approach with the standard $\mathcal{H}$-matrix approximation via ACA (labeled $\mathcal{H}-\mathrm{ACA}$ ). In the sequel, we focus on the approximation of the single-layer operator $\mathfrak{V}$. Since we assumed $\kappa h$ to be constant, we increase $\kappa$ with growing number of degrees of freedom $N$. By "acc." we label the relative error to the full matrix in the Frobenius norm.
Table 1 shows the memory consumption of the approximation of the discretization of $\mathfrak{V}$ with piecewise constant ansatz and test functions on the prolate spheroid, i.e. an ellipsoid with ten times

|  | $N$ | 6496 | 28108 | 114258 | 469010 | 1905242 |
| ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| $\kappa$ | 16 | 32 | 64 | 128 | 256 |  |
|  | mem. $[\mathrm{MB}]$ | 48 | 198 | 988 | 4754 | 20077 |
| dir $\mathcal{H}^{2}$-ACA | compr. $[\%]$ | 15 | 3.3 | 0.99 | 0.28 | 0.07 |
|  | $\mathrm{~KB} / N$ | 7.8 | 7.4 | 9.1 | 10.6 | 11.0 |
|  | time [s] | 14 | 66 | 299 | 1372 | 6079 |
|  | acc. | $5.9_{-4}$ | $1.3_{-4}$ | $1.5_{-4}$ | - | - |
|  | mem. $[\mathrm{MB}]$ | 47 | 223 | 1327 | 10432 | - |
| $\mathcal{H}$-ACA | compr. $[\%]$ | 14.7 | 3.7 | 1.33 | 0.62 | - |
|  | $\mathrm{KB} / N$ | 7.7 | 8.3 | 12.2 | 23.3 | - |
|  | time [s] | 12 | 66 | 404 | 4803 | - |
|  | acc. | $7.5_{-3}$ | $5.8_{-3}$ | $7.0_{-2}$ | - | - |

Table 1: Prolate spheroid: $\kappa h \sim 0.15, \eta_{\text {high }}=5$.
the extension in $x$-direction. The gain in both time and memory becomes visible for larger $N$. For small $N$, both methods have about the same performance. This is due to fact that directional $\mathcal{H}^{2}$ matrices adapt to the wave number and fade to usual $\mathcal{H}$-matrices for low frequencies. Due to a limited amount of storage, the largest example could be computed only for the new approach.
Observe that the achieved error of $\mathcal{H}$-ACA deteriorates for a growing number of waves, whereas $\operatorname{dir} \mathcal{H}^{2}$-ACA is able to achieve the prescribed accuracy. The reason for this behavior is that the standard matrix partitioning leads to blocks corresponding to domains on which the kernel function is asymptotically smooth with a large constant. The accuracy for $N=469010$ and $N=1905242$ was not computed, because the calculation of the difference to a full matrix in Frobenius norm takes $\mathcal{O}\left(N^{2}\right)$ time.


Figure 2: Memory (left) and time (right) on prolate spheroid for $\operatorname{dir} \mathcal{H}^{2}-\mathrm{ACA}$ and $\mathcal{H}$ - ACA .

## Neumann Problem

We consider the sound hard scattering problem and use piecewise linear ansatz and test functions for the Galerkin discretization of (32). The approximate Dirichlet datum $\left.\tilde{u}_{D} \approx u\right|_{\Gamma}$ is obtained from solving (32) with approximations $\left(\varepsilon=10^{-4}\right)$ of the discrete operators $V$ and $K$ of $\mathfrak{V}$ and $\mathfrak{K}$. We use the Neumann datum $v_{N}:=\partial_{\nu} S\left(\cdot-x_{0}\right)$ with an interior point $x_{0} \in \Omega$. In this case, we are able to compute the $L^{2}$-error $\left\|\tilde{u}_{D}-u\right\|_{L^{2}(\Gamma)} /\|u\|_{L^{2}(\Gamma)}$ to the exact Dirichlet trace given by $\left.u\right|_{\Gamma}=S\left(\cdot-x_{0}\right)$. Tables 2 and 4 show the behavior of the error on the sphere with radius 1 for fixed

| $N$ | 642 | 2562 | 10242 | 40962 | 163842 | 655362 |
| ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| $\kappa$ | 2 | 4 | 8 | 16 | 32 | 64 |
| time $[\mathrm{s}] V$ | 2 | 12 | 94 | 755 | 4571 | 27247 |
| mem. $[\mathrm{MB}] V$ | 2 | 16 | 105 | 481 | 2145 | 9409 |
| $\mathrm{~KB} / N V$ | 3.8 | 6.6 | 10.7 | 12.3 | 13.7 | 15.1 |
| time $[\mathrm{s}] K$ | 2 | 13 | 100 | 811 | 4921 | 29539 |
| mem. $[\mathrm{MB}] K$ | 2 | 16 | 105 | 471 | 2081 | 9187 |
| $L^{2}$-error | $2.6_{-3}$ | $2.1_{-3}$ | $2.0_{-3}$ | $1.9_{-3}$ | $1.9_{-3}$ | $2.0_{-3}$ |

Table 2: $\operatorname{dir} \mathcal{H}^{2}-\mathrm{ACA}: L^{2}$-error of $\tilde{u}_{D}$ with $\kappa h=0.19$.
$\kappa h$ and fixed $\kappa$, respectively. Furthermore, the CPU time and the memory consumption required by the approximations to $V$ and $K$ is shown and can be seen to behave logarithmic-linear for both fixed and growing wave numbers. As a reference, we made the same computations also using $\mathcal{H}$-matrices. The corresponding results are shown in Tables 3 and 5 . As before, the advantage of dir $\mathcal{H}^{2}$ - ACA becomes visible for a growing number of waves. It is remarkable, however, that even in the fixed frequency case the directional approach outperforms $\mathcal{H}$-ACA in terms of memory and computation time for larger degrees of freedom. For smaller degrees of freedom the performance of $\mathcal{H}$ - ACA is only slightly better.

## Visualization of Dirichlet Problem with Brakhage-Werner

We consider the sound soft scattering problem, i.e. we seek a solution $u=u_{i}+u_{s}$ of the Helmholtz equation (1), where $u_{i}(x):=\exp (\mathrm{i} \kappa(x, e))$ with $e=(1,0,1)^{T} / \sqrt{2}$ denotes the incident acoustic wave

| $N$ | 642 | 2562 | 10242 | 40962 | 163842 | 655362 |
| ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| $\kappa$ | 2 | 4 | 8 | 16 | 32 | 64 |
| time $[\mathrm{s}] V$ | 2 | 11 | 64 | 464 | 6890 | - |
| mem. $[\mathrm{MB}] V$ | 2 | 13 | 80 | 491 | 3150 | - |
| $\mathrm{KB} / N V$ | 3.1 | 5.2 | 8.2 | 12.6 | 20.2 | - |
| time $[\mathrm{s}] K$ | 3 | 14 | 77 | 520 | 7050 | - |
| mem. $[\mathrm{MB}] K$ | 2 | 12 | 78 | 478 | 3073 | - |
| $L^{2}$-error | $2.7_{-3}$ | $2.0_{-3}$ | $1.9_{-3}$ | $2.0_{-3}$ | $2.7_{-3}$ | - |

Table 3: $\mathcal{H}$-ACA: $L^{2}$-error of $\tilde{u}_{D}$ with $\kappa h=0.19$.

| $N$ | 642 | 2562 | 10242 | 40962 | 163842 |
| ---: | ---: | ---: | ---: | ---: | ---: |
| time [s] $V$ | 6 | 42 | 75 | 286 | 1613 |
| mem. $[\mathrm{MB}] V$ | 3 | 15 | 66 | 345 | 1613 |
| $\mathrm{~KB} / N V$ | 5.1 | 6.2 | 6.7 | 8.8 | 11.0 |
| time $[\mathrm{s}] K$ | 7 | 47 | 84 | 317 | 1734 |
| mem. $[\mathrm{MB}] K$ | 3 | 15 | 64 | 334 | 1669 |
| $L^{2}$-error | $1.8_{-2}$ | $4.5_{-3}$ | $1.2_{-3}$ | $3.9_{-4}$ | $2.3_{-4}$ |

Table 4: $\operatorname{dir} \mathcal{H}^{2}-\mathrm{ACA}: L^{2}$-error of $\tilde{u}_{D}$ with fixed $\kappa=6$.
and $u_{s}$ is the unknown scattered field. The incident wave is reflected on a sound soft obstacle $\Omega$, which is described by the Dirichlet condition $\left.u_{s}\right|_{\Gamma}=-\left.u_{i}\right|_{\Gamma}$.

The obstacle is composed of 4 cylindrical spheres with 507904 panels and 253960 vertices. We solve (34) for the unknown density $\phi$ with piecewise linear ansatz and test functions. Following [19], we use the coupling parameter $\beta=\kappa / 2$. In a second step, we evaluate the potential (33) on a uniform discretization of a screen behind the obstacle in order to compute the scattered wave $u_{s}$. Figure 3 shows the resulting pressure field of the total wave $\left|u_{i}+u_{s}\right|$ for $\kappa=10$ and $\kappa=40$.


Figure 3: Pressure field $\left|u_{s}+u_{i}\right|$ for $\kappa=10$ and $\kappa=40$.

## References

[1] B. K. Alpert, G. Beylkin, R. Coifman, and V. Rokhlin. Wavelet-like bases for the fast solution of second-kind integral equations. SIAM J. Sci. Comput., 14:159-184, 1993.

| $N$ | 642 | 2562 | 10242 | 40962 | 163842 |
| ---: | ---: | ---: | ---: | ---: | ---: |
| time [s] $V$ | 2 | 13 | 69 | 379 | 2771 |
| mem. $[\mathrm{MB}] V$ | 3 | 17 | 90 | 454 | 2167 |
| $\mathrm{~KB} / N V$ | 4.5 | 6.9 | 9.2 | 11.6 | 13.9 |
| time $[\mathrm{s}] K$ | 3 | 17 | 83 | 439 | 3024 |
| mem. $[\mathrm{MB}] K$ | 3 | 17 | 89 | 446 | 2122 |
| $L^{2}$-error | $1.8_{-2}$ | $4.4_{-3}$ | $1.1_{-3}$ | $3.0_{-4}$ | $1.2_{-4}$ |

Table 5: $\mathcal{H}$-ACA: $L^{2}$-error of $\tilde{u}_{D}$ with fixed $\kappa=6$.
[2] S. Amini and A. Profit. Multi-level fast multipole solution of the scattering problem. Engineering Analysis with Boundary Elements, 27(5):547-654, 2003.
[3] Christopher R. Anderson. An implementation of the fast multipole method without multipoles. SIAM J. Sci. Statist. Comput., 13(4):923-947, 1992.
[4] Lehel Banjai and Wolfgang Hackbusch. Hierarchical matrix techniques for low- and highfrequency Helmholtz problems. IMA J. Numer. Anal., 28(1):46-79, 2008.
[5] J. Barnes and P. Hut. A hierarchical $\mathcal{O}(N \log N)$ force calculation algorithm. Nature, 324:446-449, 1986.
[6] M. Bebendorf. Approximation of boundary element matrices. Numer. Math., 86(4):565-589, 2000.
[7] M. Bebendorf. Hierarchical Matrices: A Means to Efficiently Solve Elliptic Boundary Value Problems, volume 63 of Lecture Notes in Computational Science and Engineering (LNCSE). Springer, 2008. ISBN 978-3-540-77146-3.
[8] M. Bebendorf and R. Venn. Constructing nested bases approximations from the entries of non-local operators. Numer. Math., 121(4):609-635, 2012.
[9] S. Börm, M. Löhndorf, and J. M. Melenk. Approximation of integral operators by variable-order interpolation. Numer. Math., 99(4):605-643, 2005.
[10] H. Brakhage and P. Werner. Über das Dirichletsche Außenraumproblem für die Helmholtzsche Schwingungsgleichung. Arch. Math., 16:325-329, 1965.
[11] A. Brandt. Multilevel computations of integral transforms and particle interactions with oscillatory kernels. Comput. Phys. Comm., 65:24-38, 1991.
[12] A. Buffa and R. Hiptmair. A coercive combined field integral equation for electromagnetic scattering. SIAM J. Numer. Anal., 42(2):621-640, 2003.
[13] A. J. Burton and G. F. Miller. The application of integral equation methods to the numerical solution of boundary value problems. Proc. Roy. Soc., London, A232:201-210, 1971.
[14] Simon N. Chandler-Wilde, Ivan G. Graham, Stephen Langdon, and Euan A. Spence. Numericalasymptotic boundary integral methods in high-frequency acoustic scattering. Acta Numerica, 21:89-305, 2012.
[15] Hongwei Cheng, William Y. Crutchfield, Zydrunas Gimbutas, Leslie F. Greengard, J. Frank Ethridge, Jingfang Huang, Vladimir Rokhlin, Norman Yarvin, and Junsheng Zhao. A wideband fast multipole method for the Helmholtz equation in three dimensions. J. Comput. Phys., 216(1):300-325, 2006.
[16] E. Darve. The fast multipole method: numerical implementation. J. Comput. Phys., 160(1):195240, 2000.
[17] B. Engquist and L. Ying. Fast directional multilevel algorithms for oscillatory kernels. SIAM J. Sci. Comput., 29(4):1710-1737 (electronic), 2007.
[18] B. Engquist and L. Ying. A fast directional algorithm for high frequency acoustic scattering in two dimensions. Commun. Math. Sci., 7(2):327-345, 2009.
[19] K. Giebermann. Schnelle Summationsverfahren zur numerischen Lösung von Integralgleichungen für Streuprobleme im $\mathbb{R}^{3}$. PhD thesis, Universität Karlsruhe, 1997.
[20] L. Grasedyck and W. Hackbusch. Construction and arithmetics of $\mathcal{H}$-matrices. Computing, 70:295-334, 2003.
[21] L. Greengard. The rapid evaluation of potential fields in particle systems. MIT Press, Cambridge, MA, 1988.
[22] L. F. Greengard and V. Rokhlin. A fast algorithm for particle simulations. J. Comput. Phys., 73(2):325-348, 1987.
[23] L. F. Greengard and V. Rokhlin. A new version of the fast multipole method for the Laplace equation in three dimensions. In Acta numerica, 1997, volume 6 of Acta Numer., pages 229-269. Cambridge Univ. Press, Cambridge, 1997.
[24] W. Hackbusch. A sparse matrix arithmetic based on $\mathcal{H}$-matrices. Part I: Introduction to $\mathcal{H}$ matrices. Computing, 62(2):89-108, 1999.
[25] W. Hackbusch and B. N. Khoromskij. A sparse $\mathcal{H}$-matrix arithmetic: general complexity estimates. J. Comput. Appl. Math., 125(1-2):479-501, 2000. Numerical analysis 2000, Vol. VI, Ordinary differential equations and integral equations.
[26] W. Hackbusch and B. N. Khoromskij. A sparse $\mathcal{H}$-matrix arithmetic. Part II: Application to multi-dimensional problems. Computing, 64(1):21-47, 2000.
[27] W. Hackbusch, B. N. Khoromskij, and S. A. Sauter. On $\mathcal{H}^{2}$-matrices. In H.-J. Bungartz, R. H. W. Hoppe, and Ch. Zenger, editors, Lectures on Applied Mathematics, pages 9-29. Springer-Verlag, Berlin, 2000.
[28] W. Hackbusch and Z. P. Nowak. On the fast matrix multiplication in the boundary element method by panel clustering. Numer. Math., 54(4):463-491, 1989.
[29] M. Messner, M. Schanz, and E. Darve. Fast directional multilevel summation for oscillatory kernels based on Chebyshev interpolation. J. Comput. Phys., 231:1175-1196, 2012.
[30] V. Rokhlin. Rapid solution of integral equations of classical potential theory. J. Comput. Phys., 60(2):187-207, 1985.
[31] V. Rokhlin. Rapid solution of integral equations of scattering theory in two dimensions. J. Comput. Phys., 86(2):414-439, 1990.
[32] V. Rokhlin. Diagonal forms of translation operators for the Helmholtz equation in three dimensions. Appl. Comput. Harmon. Anal., 1(1):82-93, 1993.
[33] E. E. Tyrtyshnikov. Mosaic-skeleton approximations. Calcolo, 33(1-2):47-57 (1998), 1996. Toeplitz matrices: structures, algorithms and applications (Cortona, 1996).
[34] Lexing Ying, George Biros, and Denis Zorin. A kernel-independent adaptive fast multipole algorithm in two and three dimensions. J. Comput. Phys., 196(2):591-626, 2004.

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[^0]:    *This work was supported by the DFG project BE2626/3-1.

